**Predicting Materials and** 10.811 Material Properties by **Density Functional Theory:** 

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BORON



Jens Kunstmann

Max-Planck-Institut für Festkörperforschung Stuttgart, Germany



- Motivation
- Boron Sheets
- Boron Nanotubes
- Structure Control
- Nanotubular Junctions
- Layered Bulk Phases of Boron
- Enatom Method
- Summary

# Motivation

### Why is elemental boron interesting?

- little studied
- many fundamental properties are unknown (phase diagram, ground state structure, etc.)
- very complex chemistry and bonding

### **Bulk Phases**

- 3D frameworks of B<sub>12</sub> icosahedra
- different bulk phases
  - $\rightarrow$   $\alpha$ -rhombohedral (12 at./cell)
  - $\rightarrow$   $\beta$ -rhombohedral (106.7 at./cell)
  - $\rightarrow$   $\beta$ -tetragonal (189.9 at./cell)
- semiconducting
- **superconducting** at high pressure Eremets *et al.*, Science **293**, 272 (2001)



α-rhomb. boron

# Motivation

### **Clusters**

• Aufbau principle: Boustani, Phys. Rev. B 55 (1997)

→ prediction of boron nanostructures: nanotubes, fullerenes, sheets

- small boron clusters are quasiplanar
  - $\rightarrow$  theory:I. Boustani, Surf. Sci. **370** (1997) $\rightarrow$  experiments:H.J. Zhai, et al. Nat. Mater. **2** (2003)







- boron nanotubes
  - $\rightarrow$  theory:
  - $\rightarrow$  experiments:

Boustani *et al.*, Europhys. Lett. **39** (1997) Ciuparu *et al.*, J. Phys. Chem. B **108** (2004)



# Motivation

#### Many open questions:

- What is the **detailed** atomic structure of the boron nanostructures?
- What are the properties of these nanostructures?
- How can we understand the chemical bonding in the different phases?
- What is the ground state structure/phases diagram?
- What bulk phase is responsible for the highpressure superconductivity?

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#### **Density Functional Theory (DFT)**

- calculate electronic structure (charge density, DOS, band structure, Fermi surface)
- structural simulations and optimizations
- Codes:

   → TB-LMTO-ASA: tight-binding LMTO in atomic sphere approximation
   → LmtART: all-electron full-potential LMTO
   → VASP: plane waves and pseudopotentials

### **Boron Sheet**

#### **Experiment:**

- nanotubes Ciuparu *et al.*, J. Phys. Chem. B **108** (2004)
- small quasiplanar clusters H.J. Zhai, et al. Nat. Mater. 2 (2003)





#### **Broad Boron Sheet:**

- the precursor of boron nanotubes
- the boron analogue of a single graphite sheet

# What does a broad boron sheet look like?

### **Boron Sheet**

top view



side view

triangular lattice

simple up-and-down structure

- structure found by DFT simulations
- structure independently found also by Evans *et al.*, Phys. Rev. B **72**, 45434 (2005). Cabria *et al.*, Nanotechnology **17**, 778 (2006).

JK et al., Phys. Rev. B 74 (2006)

### **Boron Sheet**



yellow: charge density contour at 0.9 e/Å<sup>3</sup>

- anisotropic in-plane mechanical properties
- metallic properties

**Boron Nanotubes:** 

All boron nanotubes are metallic!

### **Boron Nanotubes**



zigzag

### armchair BNT:

 bent circumferential σ bonds



- straight axial  $\sigma$  bonds
- NO circumferential σ bonds

### **Boron Nanotubes**





 unique property among nanotubular materials usually E<sub>strain</sub> = E<sub>strain</sub>(R) only



#### suggestion:

- control radius R by growing nanotubes out of porous catalysts with well defined pores sizes (e.g. Mg-MCM-41)
- control chiral angle θ by tuning the reaction conditions
   synthesis of a specific (R,θ) nanotube



Anisotropic in-plane mechanical properties of sheets allow for more structure control of the tubes!

### **Boron Nanotubes**

- Summary: theory of boron nanotubes (BNTs): unifying and generalizing former studies in the field JK *et al.*, Chem Phys. Lett. **402** (2005). JK *et al.*, Phys. Rev. B **74** (2006).
  - all BNTs are metallic
  - the strain energy of a BNT is a function of the tube's radius and chiral angle:  $E_{strain} = E_{strain}(R,\theta)$
  - new approach to the long standing problem of structure control in nanoscience JK et al., Nanotechnology 18, 155703 (2007).

#### further developments:

revised version of our theory:

Tang *et al.*, Phys. Rev. Lett. **99** (2007). Yang *et al.*, Phys. Rev. B **77** (2008).

- There are many *different* nanotubular materials: existing: BN, MoS<sub>2</sub>, VO<sub>x</sub>, ... predicted: MB<sub>2</sub>, B<sub>2</sub>O, CaSi<sub>2</sub>, ...
- Idea:

Heterogeneous Nanotubular Networks

### • Basic question:

Is it possible to *connect* different nanotubular materials?

#### • Structural paradigm:

intamoledular junctions between boron and carbon nanotubes

**Model System:** linear BC-heterojunction of arbitrary radius and chirality from a single sheet



• Only structure elements | and II can appear at the interface

#### **Results:**

DFT simulations of different model junctions that only contain structure element



#### Solution:

• Linear BC-heterojunctions of any type can be formed by incorporating structure element **II** at the interface.



Heterogeneous Nanotubular Networks are possible!

#### Theory:

- electrons, phonons, and EP interaction are described in reciprocal space
- we want a real-space desription of
  - $\rightarrow$  solids
  - $\rightarrow$  vibrational properties
  - $\rightarrow$  electron-phonon interactions

### Enatom:

- introduced by M. A. Ball as generalized pseudoatom J. Phys. C 8, 3328 (1975)
- from greek:

en = inside atom = the indivisible part enatom = an atom inside a solid

JK et al., Phys. Rev. B **75** (2007)

**Definition:** 

- atomic displacement:  $\delta R_j = R_j R_j^o$
- definition via Helmholz decomposition of first order change in charge density:

$$rac{\partial n(oldsymbol{r})}{\partial oldsymbol{R}_j} = -
abla 
ho_j(oldsymbol{r}-oldsymbol{R}_j^o) + 
abla imes oldsymbol{B}_j(oldsymbol{r}-oldsymbol{R}_j^o)$$

- $n(\mathbf{r})$  ... total charge density
- $ho_j(\mathbf{r})$  ... rigid density (scalar field); describes the charge density that moves rigidly with the atom upon a displacement
- $abla imes B_j(r)$  ... deformation density (vector field); describes how the charge density deforms due to a nuclear displacement

Properties:

• 
$$\sum_{j} \rho_j(\boldsymbol{r} - \boldsymbol{R}_j^o) = n(\boldsymbol{r})$$

The enatom is a unique decomposition of the total charge density into a sum of atomic-like contributions.

• charge density of displaced atoms:

$$n(\boldsymbol{r}; \{\boldsymbol{R}_j\}) = \sum_j [\rho_j(\boldsymbol{r} - \boldsymbol{R}_j^o - \delta \boldsymbol{R}_j) + \delta \boldsymbol{R}_j \cdot \nabla \times \boldsymbol{B}_j(\boldsymbol{r} - \boldsymbol{R}_j^o)]$$

This holds for a first order displacement of the atoms from their equilibrium position.



The same construction can also be applied to the total **potential** 

$$egin{array}{rcl} n(m{r}) & o & v(m{r}) \ 
ho_j(m{r}) & o & V_j(m{r}) \ 
abla imes m{B}_j(m{r}) & o & 
abla imes m{W}_j(m{r}) \end{array}$$

#### **Example:** enatom density of fcc Li at P=35 GPa

#### rigid density ρ





countour plot of  $\boldsymbol{\rho}$  in xy plane

#### **Example:** enatom density of fcc Li at P=35 GPa

#### deformation density $\nabla \times \mathbf{B}$





countour plot of  $|\nabla \mathbf{x} \mathbf{B}|$  in xy plane

#### **Results:**

- first realization of the enatom method
- studied pressure evolution of enatom quantities for fcc Li and fcc Al

### Outlook:

- generalization of the enatom code for more complex systems
- study of non-metals
- calculation of phonon frequencies
- calculation of electron-phonon matrix elements
- study boron



• generalized theory for boron nanotubes

JK *et al.*, Chem Phys. Lett. **402** (2005). JK *et al.*, Phys. Rev. B **74** (2006).

 new route for structure control of nanotubes during synthesis

JK et al., Nanotechnology 18, 155703 (2007).

- intramolecular junctions of nanotubes JK *et al.*, J. Chem. Phys. **121** (2004)
- proposed layered bulk phases of boron JK *et al.*, (to be published)
- first realization of entaom method JK et al., Phys. Rev. B **75** (2007)

### **Thanks**

- Ole K. Andersen (MPI Stuttgart)
- Lilia Boeri (MPI Stuttgart)
- Alexander Quandt (Uni Greifswald)
- Warren E. Pickett (UC Davis, USA)
- Jens Kortus (TU Freiberg)
- Ihsan Boustani (Uni Wuppertal)