### **DNA Conductivity: Our Most Recent Results**

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# MOTIVATION



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### **Topics for today:**

- DNA conductivity is <u>temperature</u> <u>independent</u>

- <u>Acoustic modes</u> in DNA duplexes and their role in DNA charge transport **DNA conductivity is** <u>*temperature independent*</u> <u>**Two experimental techniques (Richard Nichols** <u>*et al.*</u>)</u>



a) shows the *I*(s)-technique where the GOLDEN tip is withdrawn from the GOLDEN surface while a DNA molecule bridges the gap.
b) shows the *I*(*t*)-technique where the GOLDEN tip - GOLDEN substrate gap is fixed and a DNA molecule spontaneously bridges the gap.



Characteristic results for both *I*(*s*) and *I*(*t*)-technique are shown in c) and d) respectively.

The resulting histograms are shown in e) (*I*(*s*)) and f) (*I*(*t*)).



The jump current using the *I(t)* method as a function of the tip substrate potential for both (dA)<sub>15</sub>-(dT)<sub>15</sub> and (dG)<sub>15</sub>-(dC)<sub>15</sub>.



The single molecule conductance as a function of the substrate temperature at zero bias for both (dA)<sub>15</sub>-(dT)<sub>15</sub> and (dG)<sub>15</sub>-(dC)<sub>15</sub>.

DNA duplex	Conductance / nS
HS-C <sub>3</sub> -(dA)	$0.51 \pm 0.08$
(dT)-C <sub>3</sub> -SH	0.51 ± 0.08
HS-C <sub>3</sub> -(dG)	12+02
(dC)-C <sub>3</sub> -SH	$1.2 \pm 0.2$

The jump current using the *I(t)* method as a function of the tip substrate potential for both (dA)<sub>15</sub>-(dT)<sub>15</sub> and (dG)<sub>15</sub>-(dC)<sub>15</sub>.

### <u>Acoustic modes</u> in DNA duplexes <u>Internal energy for the Debye model of a crystal lattice</u> <u>with 3 acoustic and 3s-3 optical modes</u>

$$\mathcal{E} = \mathcal{E}_0 + Nk_B T \left[ 3\mathcal{D} \left( \frac{\mathcal{O}}{T} \right) + \sum_{j=4}^{3s} \frac{\mathcal{O}_j / T}{e^{\mathcal{O}_j / T} - 1} \right]$$

where the *Debye function*  $\mathcal{D}(z)$  is

$$\mathcal{D}(z) = \frac{3}{z^3} \int_0^z \frac{x^3 dx}{e^x - 1}$$

At high temperatures,  $T >> \Theta_j$  (and, consequently,  $T >> \Theta$ ) we get z << 1 and  $\mathcal{D}(0) = 1$ 

## And what is the Debye temperature for a typical DNA ???

## It is about 160° K,

if one takes the <u>velocity of sound</u> in B-DNA to be about 1900 m/sec (M. B. Hakim, S. M. Lindsay, and J. Powell, *Biopolymers*, 1984, vol. 23, p. 1185) and the B-DNA *lattice constant* of 3.4 Å.

whereas the DNA measurements were carried out at <u>room temperature</u> !!! <u>Therefore, one may speak of</u> <u>free acoustic polarons</u> <u>as possible charge carriers</u> responsible for DNA conductivity.

To check this hypothesis, we have to know,

## 1) what DNA acoustic modes are looking like and

2) whether free acoustic polarons are possible along DNA duplexes

### **DNA acoustic modes**

# can be estimated using *"elastic network" model*

## (coarse-graining model of DNA assuming that nucleotides are rigid bodies)

### **DNA acoustic mode 1**



If Z is the double-helical axis and the ends are fixed by electrodes,



it is "guitar string" <u>BENDING</u> in X-direction.

### **DNA acoustic mode 2**



If Z is the doublehelical axis and the ends are fixed by electrodes,



This is "guitar string" <u>BENDING</u> in Y-direction.

### **DNA acoustic mode 3**





This mode is much more than just the 2nd harmonics of the "guitar string" <u>BENDING,</u> since DNA is compressible:

It is *TWISTING* together with changes in the *double-helical PITCH*  <u>To check numerically, whether free</u> <u>acoustic polarons are possible in DNA</u> <u>duplexes:</u>

1) to generate deformed DNA conformations according to its acoustic modes

2) to calculate and compare conductance along DNA duplexes in their equilibrium and deformed conformations **DNA static ballistic conductivity** 

101

$$G = \lim_{\omega \to 0} \frac{\langle I \rangle}{V} = -e^2 \hbar \pi \lim_{\omega \to 0} \sum_{\alpha, \beta} \left| \langle \alpha | \hat{v}_x | \beta \rangle \right|^2 \frac{f_\alpha - f_\beta}{\varepsilon_\alpha - \varepsilon_\beta} \,\delta \left( \varepsilon_\alpha - \varepsilon_\beta - \hbar \omega \right)$$

<u>Above is the</u> <u>conventional Kubo formula,</u> <u>in terms of eigenstates and eigenfunctions</u>

$$G = 2\left(\frac{e^{2}}{h}\right) \operatorname{Tr}\left[(i\hbar\hat{v}_{x})\operatorname{Im}\hat{\mathcal{G}}(E)(i\hbar\hat{v}_{x})\operatorname{Im}\hat{\mathcal{G}}(E)\right] \qquad \frac{...\ recast}{in\ the}$$
$$\operatorname{Im}\hat{\mathcal{G}}(E) = \frac{1}{2i}\left[\hat{\mathcal{G}}^{R}(E) - \hat{\mathcal{G}}^{A}(E)\right] \qquad \frac{Green}{function}$$

*To estimate DNA conductivity we* <u>need:</u>

# • <u>Velocity of the charge</u> <u>flow</u> through DNA

• <u>Green function</u> of DNA (in the standard <u>Extended Hückel</u> approximation)





$$V = \theta$$

$$I = 0$$

$$\left[E\mathbf{I} - \mathbf{H} - \boldsymbol{\Sigma}_{l}(E) - \boldsymbol{\Sigma}_{r}(E)\right]\hat{\mathcal{G}}(E) = \mathbf{I}$$

2. Full Green function of the contact between the left lead and the 1st layer



The Fermi level (defined here as halfway between the energies of the highest occupied and lowest unoccupied orbitals at 0 K) is around -12.05 eV for all graphs. Red graphs correspond to the transmission spectra of the duplexes in their ideal B-DNA conformation, the green ones depict those in the presence of deformations along DNA acoustic modes

**Conclusions** 

Acoustic polaron formation promotes charge transport in (dA)-(dT) systems

Acoustic polaron formation interfers charge transport in (dG)-(dC)

**Outlook** 

## Influence of DNA dynamics on the faster time scales

The exact molecular structure of transmissive pathways (conductive channels) in DNA

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