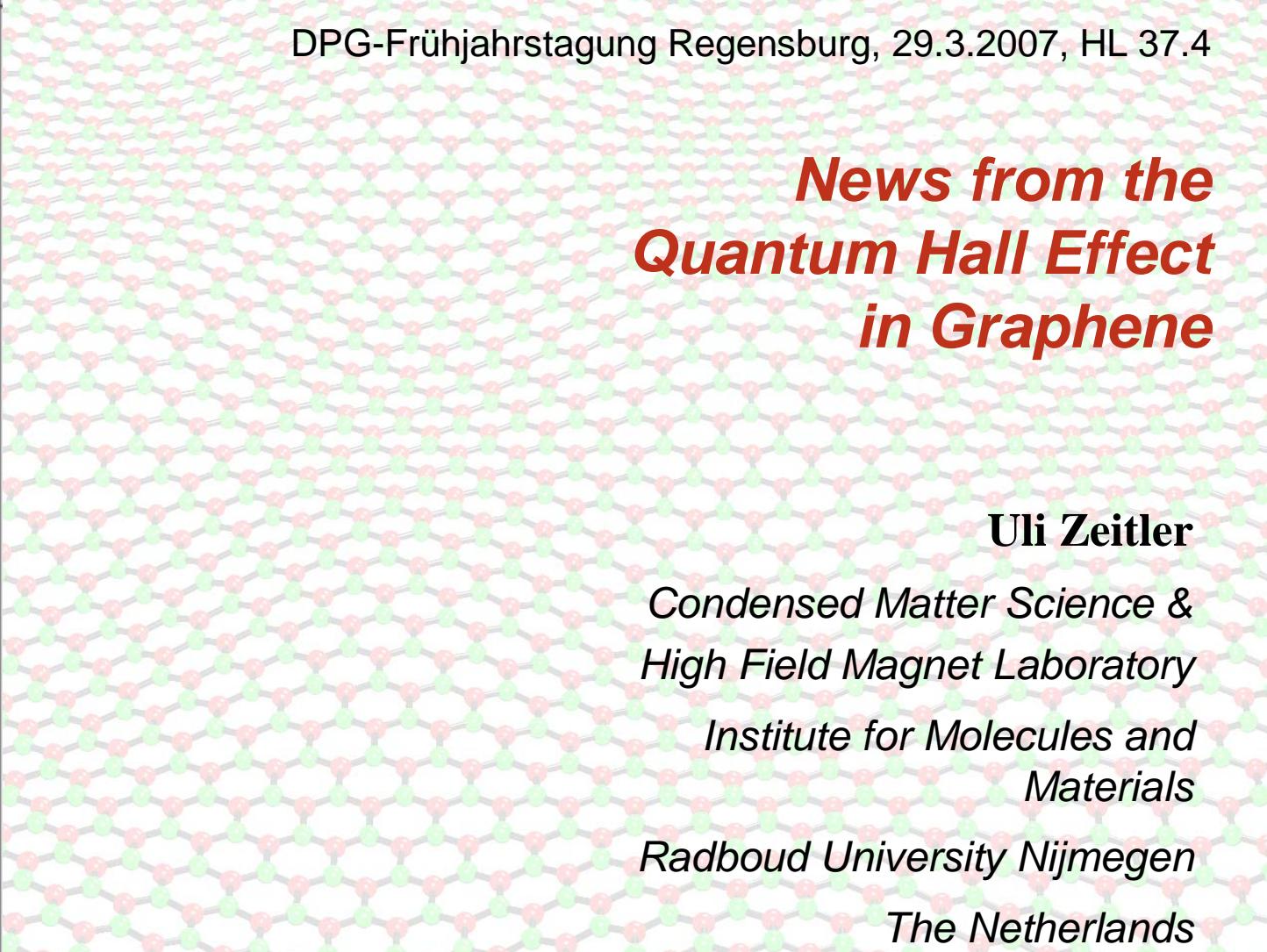


FACULTY OF SCIENCE, MATHEMATICS AND COMPUTING SCIENCE

DPG-Frühjahrstagung Regensburg, 29.3.2007, HL 37.4

A repeating pattern of a hexagonal lattice structure, representing graphene, composed of red and green spheres connected by grey lines.

***News from the  
Quantum Hall Effect  
in Graphene***

**Uli Zeitler**

*Condensed Matter Science &  
High Field Magnet Laboratory*

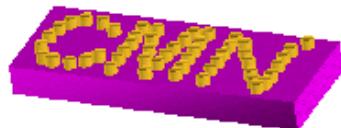
*Institute for Molecules and  
Materials*

*Radboud University Nijmegen  
The Netherlands*



QHE in Graphene

## Collaborators



Kostya Novoselov, S.V. Morozov,\* D. Jiang,  
F. Schedin, and Andre Geim

MANCHESTER  
1824

*Manchester Centre for Mesoscience and Nanotechnology,  
University of Manchester. UK*

\**Institute for Microelectronics Technology, Chernogolovka, Russia*

Misha Katsnelson

*Condensed Matter Theory, IMM, Radboud University Nijmegen*



Jos Giesbers, Jan Kees Maan, and UZ

*Condensed Matter Science / High Field Magnet Laboratory  
Institute for Molecules and Materials, Radboud University Nijmegen*



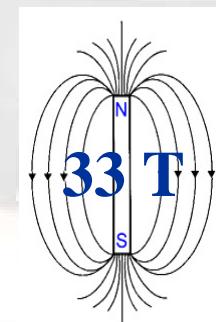
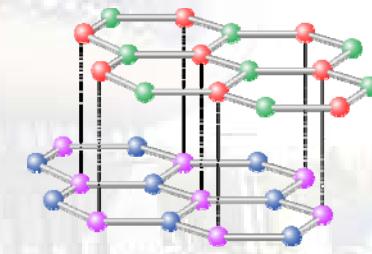
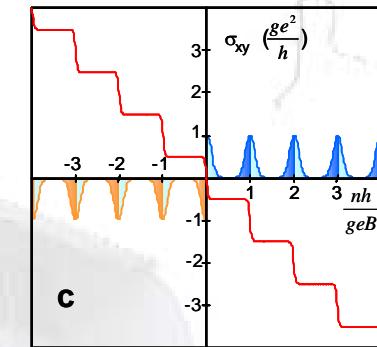
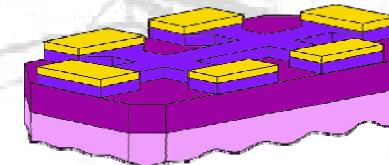
DPG-Regensburg, 29.3.2007

HFML

*Science in High Magnetic Fields*



- Introduction
  - Graphene devices
  - Quantum Hall effect
- Single-layer graphene
  - Bandstructure of graphene
  - Half-integer QHE
  - Room-temperature QHE
- Bilayer graphene
  - Bandstructure of massive chiral Dirac fermions
  - A new type of integer QHE
- Conclusions
- Commercial



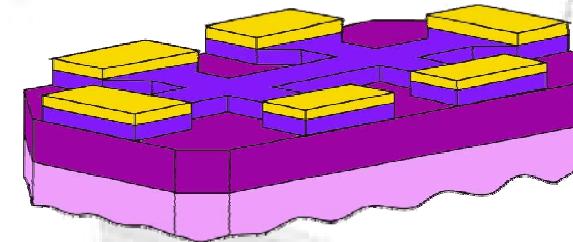


# Introduction: Graphene devices

## ➤ Introduction

### ➤ Graphene devices

- Quantum Hall effect
- Single-layer graphene
  - Bandstructure of graphene
  - Half-integer QHE
  - Room-temperature QHE & activation gaps
- Bilayer graphene
  - Bandstructure of massive chiral Dirac fermions
  - A new type of integer QHE
- Conclusions

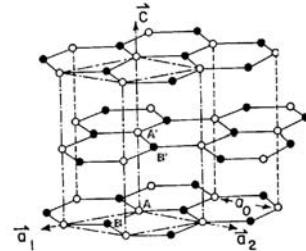




QHE in Graphene

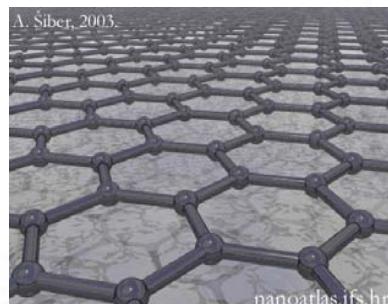
# Carbon allotropes

3D



*Graphite*

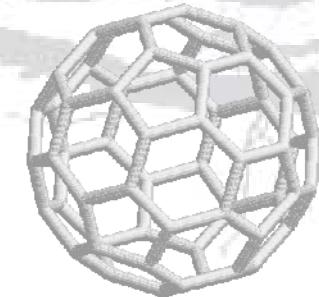
2D



*Graphene*

*Novoselov et al., 2004*

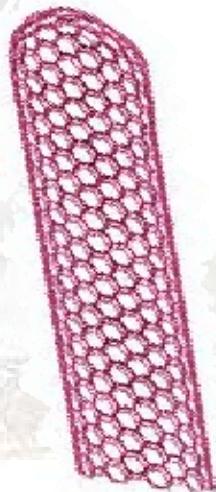
0D



*Fullerenes*

*Curl, Kroto, Smalley 1985  
Nobel prize 1996*

1D



*Carbon  
Nanotubes*

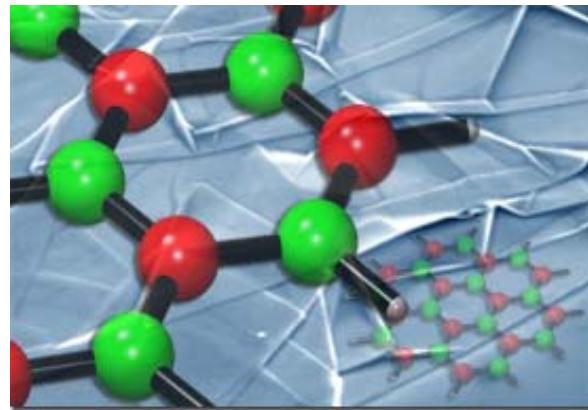
Multi-wall 1991  
Single-wall 1993



DPG-Regensburg, 29.3.2007



Science in High Magnetic Fields



- mono-atomic layer of carbon
- first truly 2D system  
(*that should not exist !!*)
- unique electronic properties  
(zero-gap semiconductor with  
**massless chiral Dirac fermions**)

MANCHESTER  
1824

2004

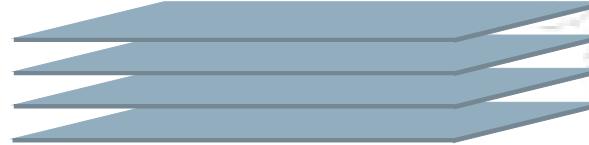
A monolayer of graphite  
on SIMOX is visible  
under an **optical microscope**

Other truly 2D systems

*Novoselov et al., PNAS 2005*

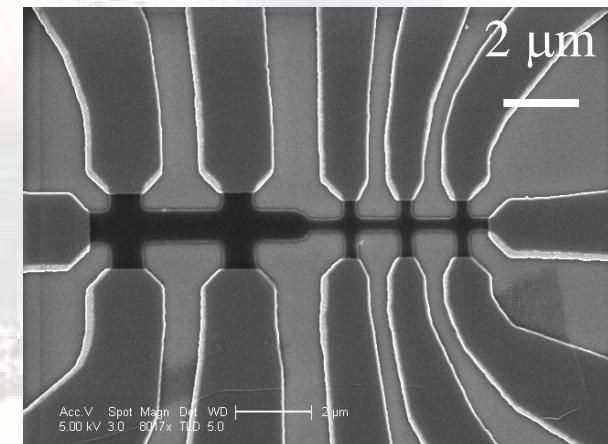
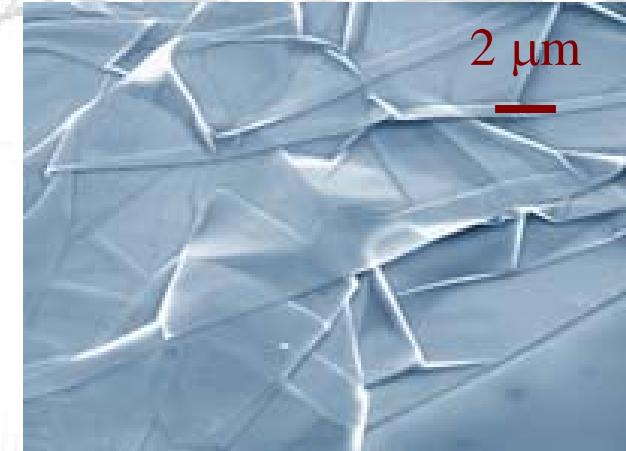


# Making graphene devices



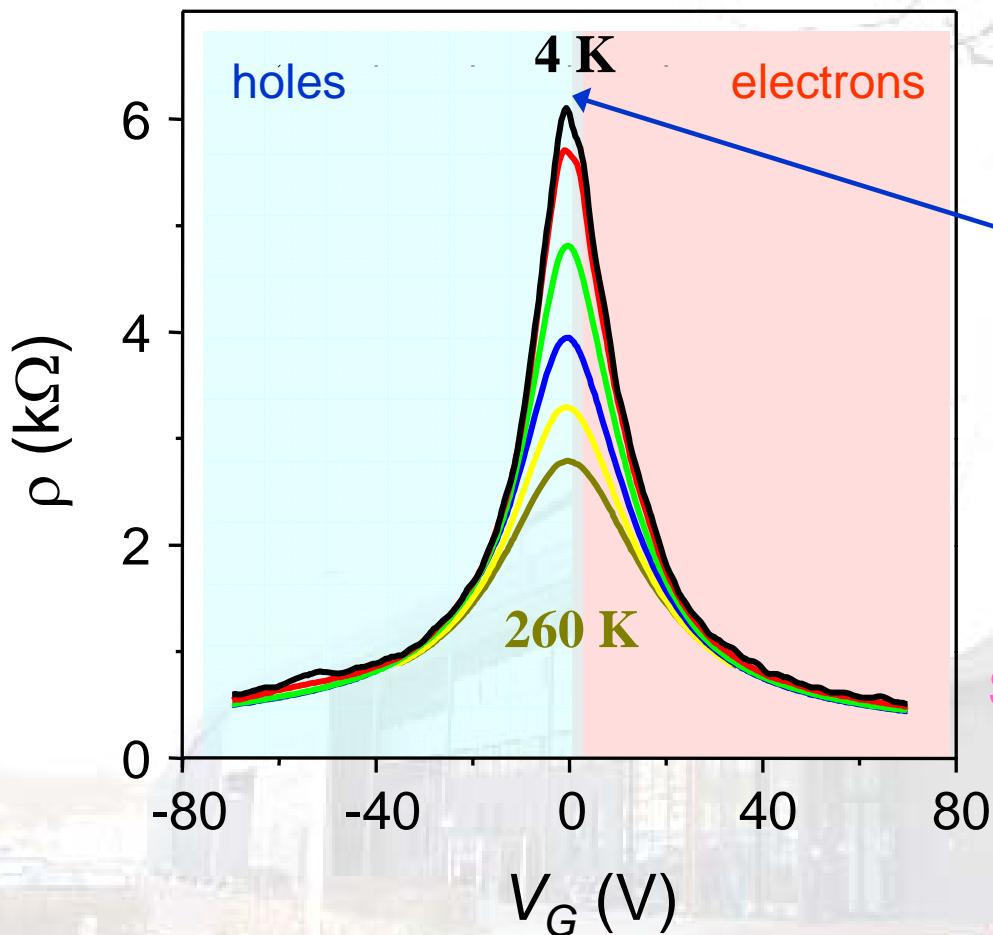
- Mechanical cleavage of natural graphite  
*“peel off down to a single layer”*
- Identification of single-layer flakes under ***optical microscope***
- Contacting with Au contacts
- Mesa etching

Novoselov *et al.*, Science 2004

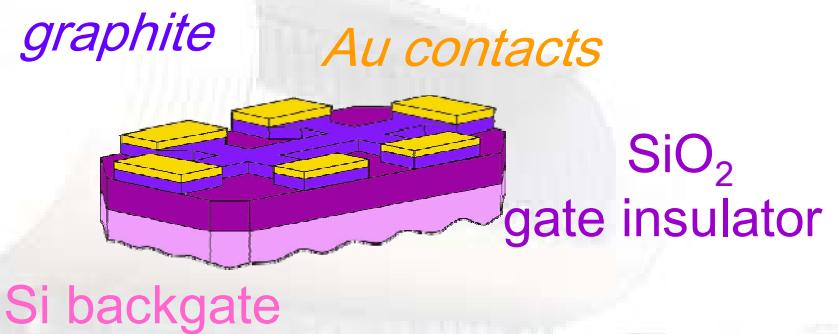




# Ambipolar field-effect transistor



Field effect doping in a MOSFET



Novoselov et al.,  
Science 2004, PNAS 2005

conducting without any charge carriers !!!



# Introduction: Quantum Hall effect

## ➤ Introduction

- Graphene devices

## ➤ Quantum Hall effect

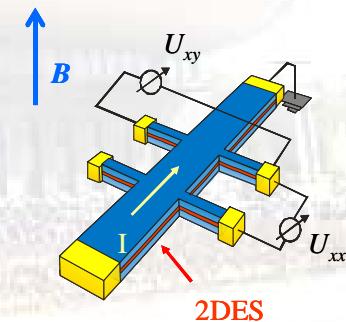
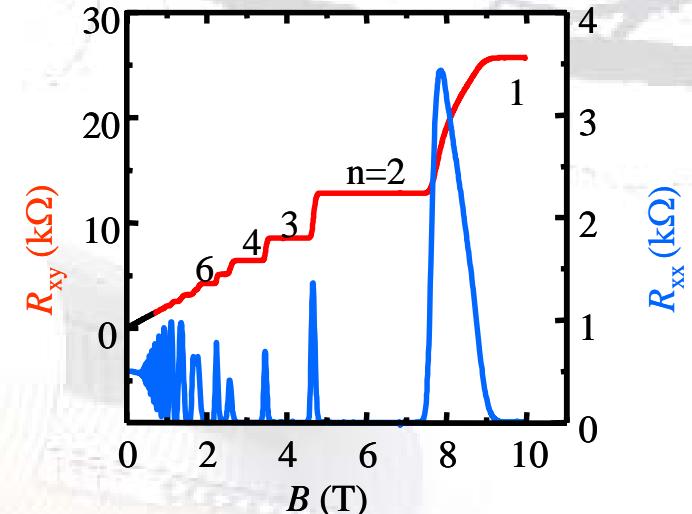
### ➤ Single-layer graphene

- Bandstructure of graphene
- Half-integer QHE
- Room-temperature QHE

### ➤ Bilayer graphene

- Bandstructure of massive chiral Dirac fermions
- A new type of integer

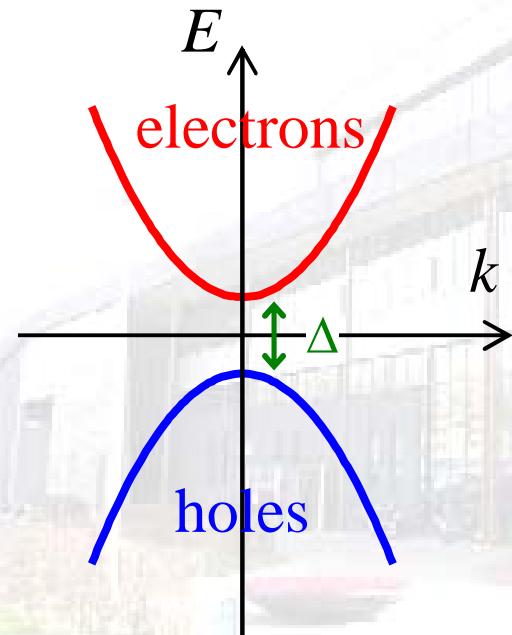
### ➤ Conclusions



## Landau levels

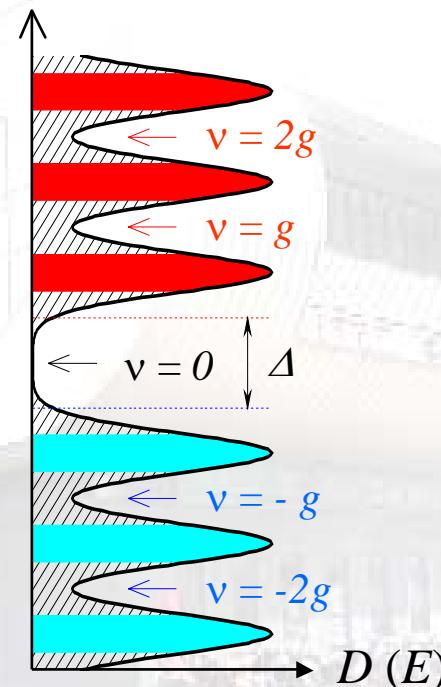
### Parabolic Dispersion

$$E^{(e/h)} = \pm \left( \frac{\hbar^2 k^2}{2m^*} + \frac{\Delta}{2} \right)$$



### Landau quantisation

$$E_N = \pm [\hbar \omega_c (N + \frac{1}{2}) + \Delta / 2]$$



Level-degeneracy:

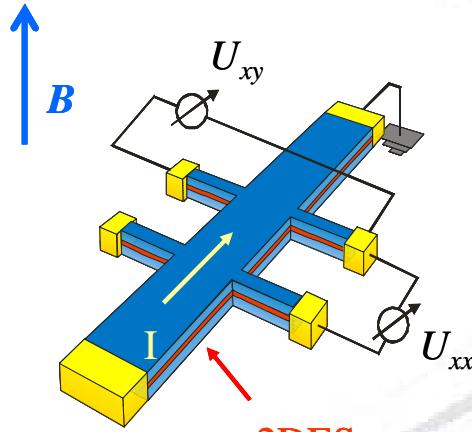
$$n_L = g_v g_s \frac{eB}{h} = 4 \frac{eB}{h}$$

*independent on  
dispersion!*

*localised* and *extended* states

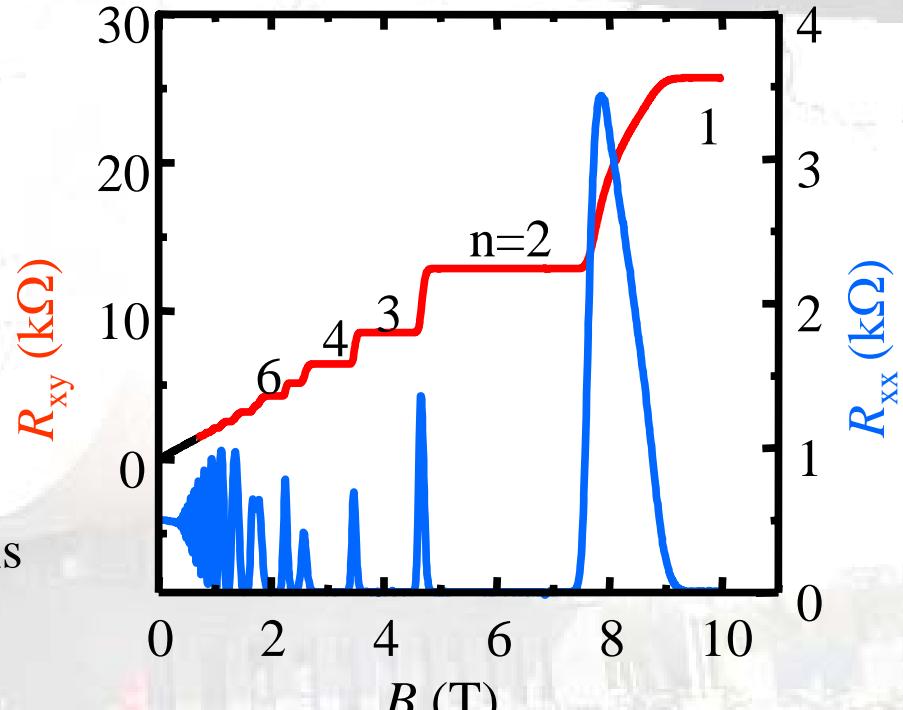


# Magnetotransport



sweep magnetic field

- ⇒ successively depopulate Landau levels
- ⇒  $1/B$ -periodic oscillations in  $R_{xx}$
- ⇒ quantized plateaus in  $R_{xy} = (h/ne^2)$
- ⇒  $R_{xx} \rightarrow 0$  in quantum Hall plateaus

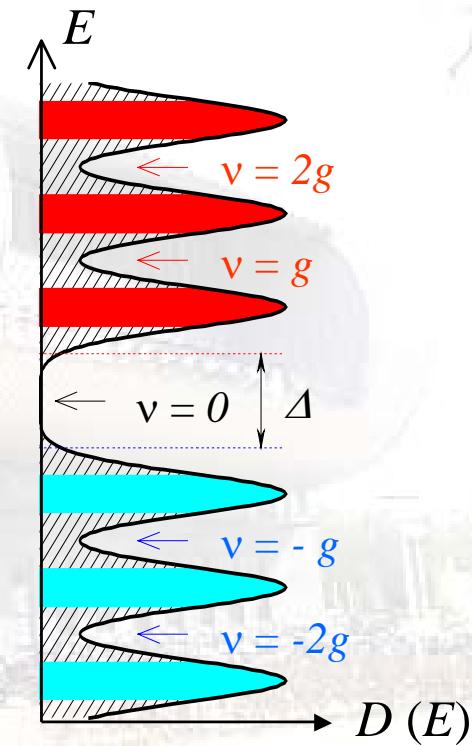
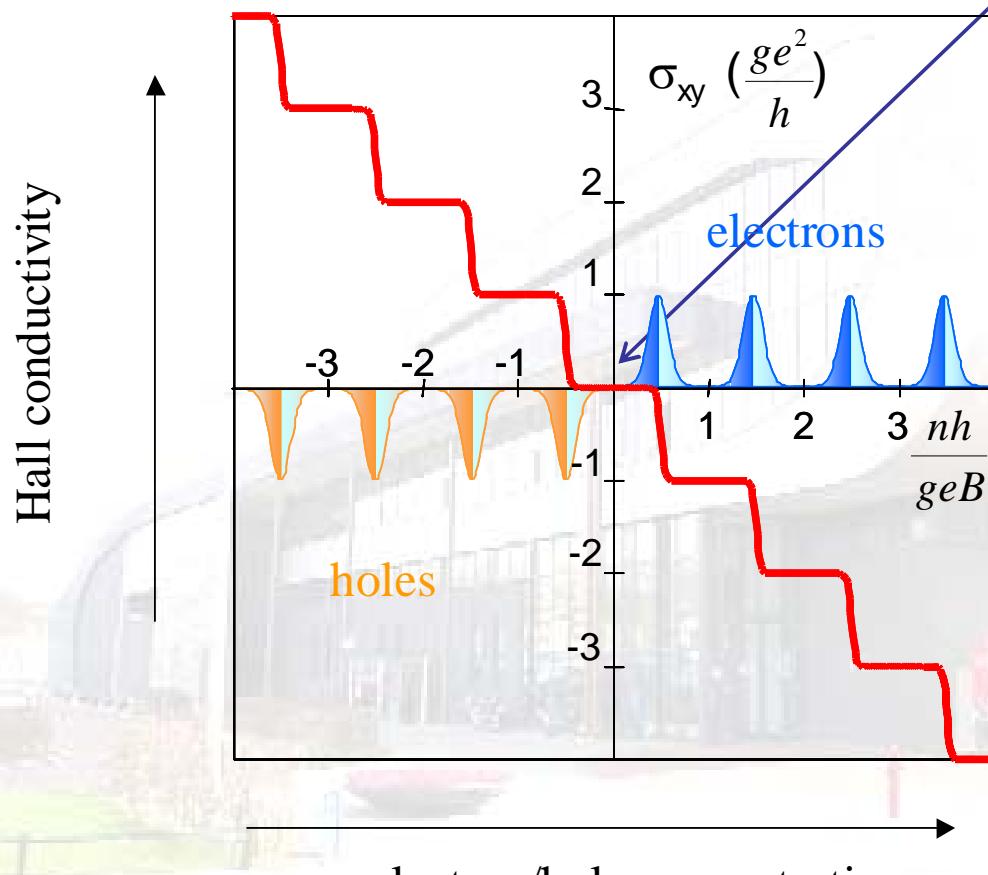




# Conventional quantum Hall effect

- quantized Hall conductivity at integer Landau level filling

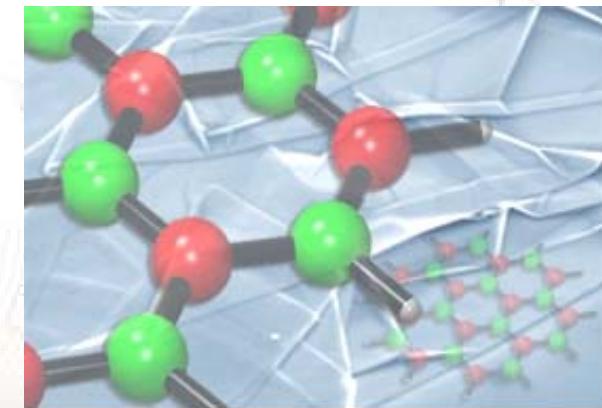
insulating at charge neutrality point





# Single-layer graphene

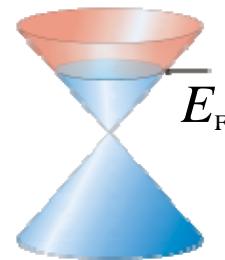
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- Conclusions





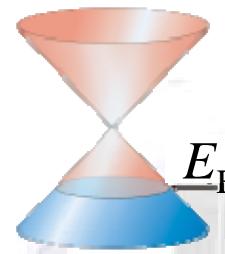
# Bandstructure graphene

*n*-doped



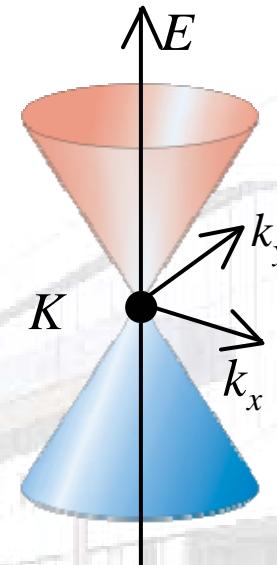
$$E(k) = \pm c \hbar |k|$$

$$E = c \hbar k \quad c \approx 10^6 \text{ m/s}$$

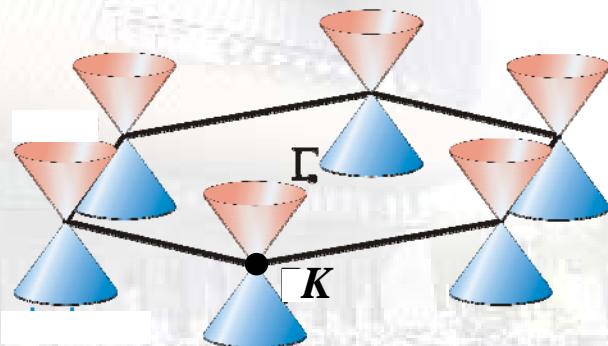


*p*-doped

Fermi energy  
exactly on Dirac point



zero-gap semiconductor  
linear dispersion  
(relativistic particles)



Wallace, 1947



# Dirac equation in a magnetic field

Schrödinger

$$\left( \frac{1}{2m^*} (\hat{\mathbf{p}} - e\mathbf{A})^2 \right) \psi = E \psi$$

$$E_N = \frac{\hbar e B}{m^*} (N + \frac{1}{2})$$

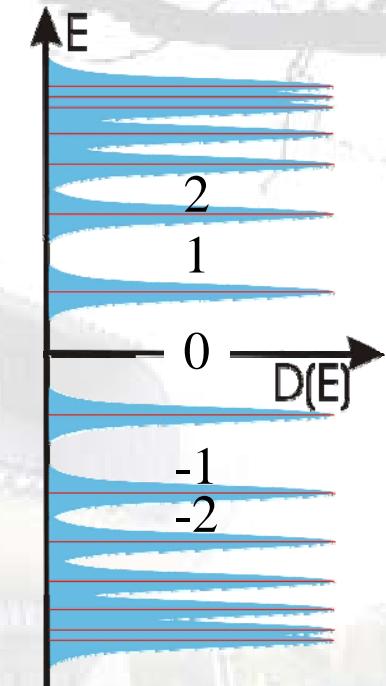
Dirac

$$(c(\hat{\mathbf{p}} - e\mathbf{A})) \psi = E \psi$$

$$(c^2 (\hat{\mathbf{p}} - e\mathbf{A})^2) \psi = E^2 \psi$$

$$E_N^2 = 2c \hbar e B (N + \frac{1}{2})$$

$$E_N = \pm \sqrt{2c \hbar e B (N + \frac{1}{2})}$$



integer QHE !

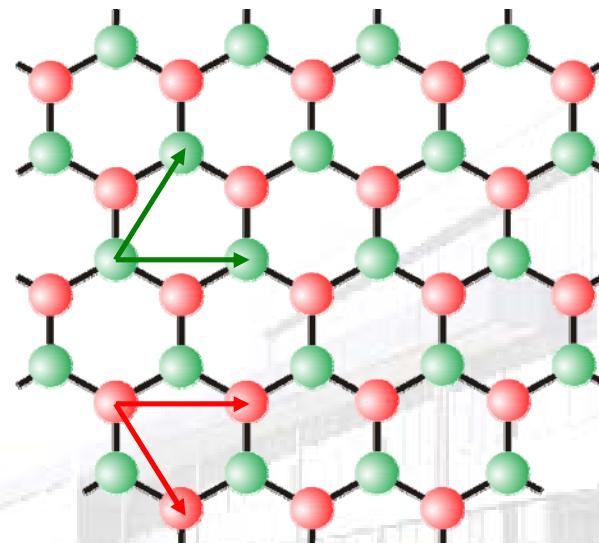


# Chirality

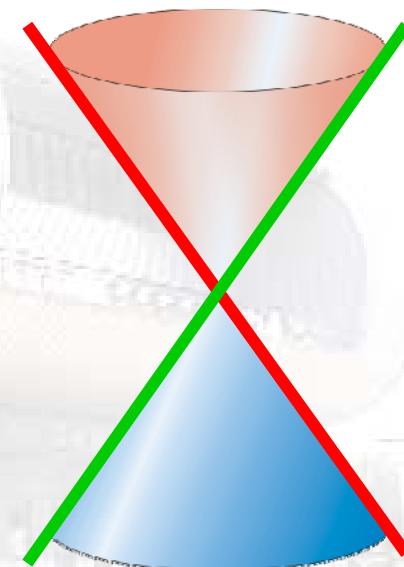
**honeycomb lattice:**

**2 triagonal sublattices ( +/ - )**

$$\psi = \sin \varphi |+\rangle + \cos \varphi |-\rangle$$



Chiral superposition of  
sublattice wave functions  
(pseudospin)



# Chiral Dirac fermions in a magnetic field

relativistic Landau levels

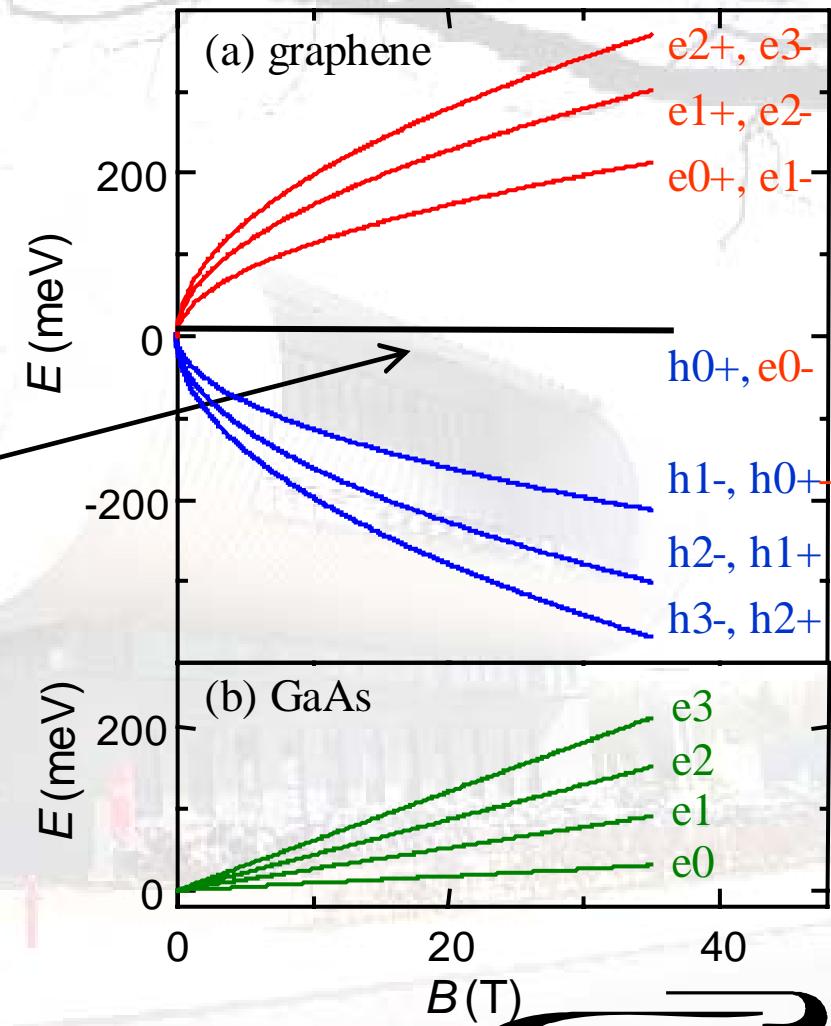
$$E_{N,\pm} = \pm \sqrt{2\hbar c^2 e B \left(N + \frac{1}{2} \pm \frac{1}{2}\right)}$$

chirality index

Landau level **at zero energy**  
*shared between electrons and holes*

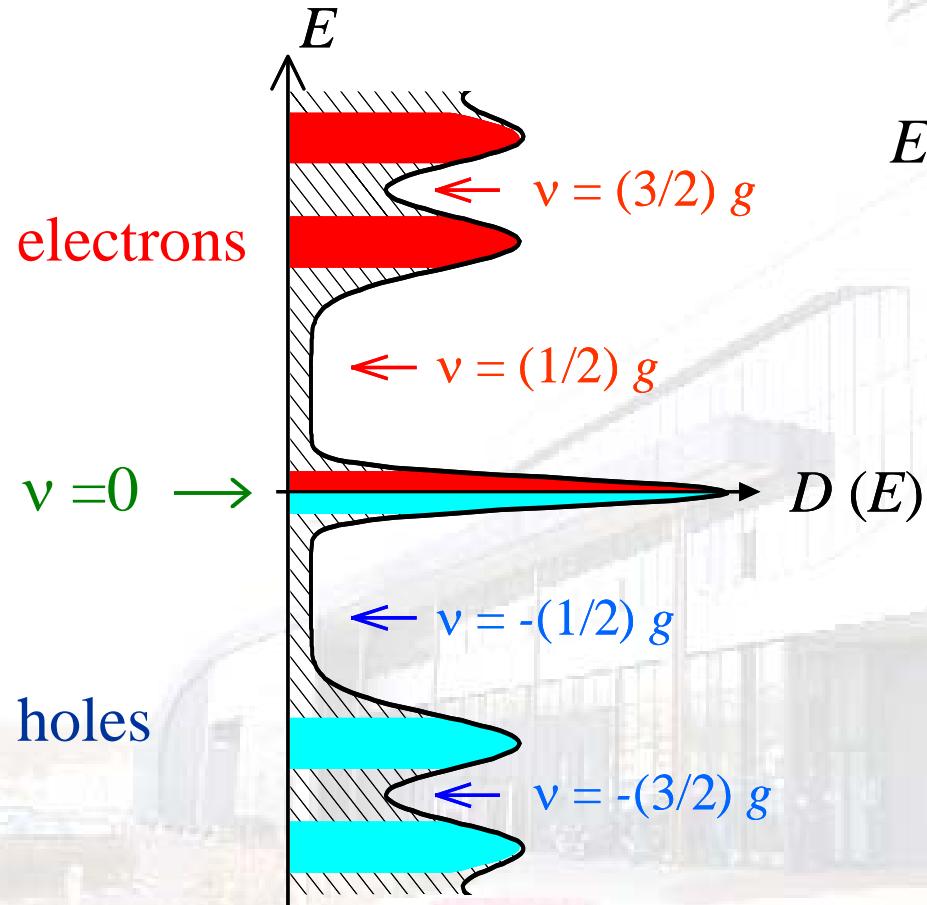
traditional Landau levels

$$E_N = \pm \hbar \omega_c \left(N + \frac{1}{2}\right)$$





# Landau levels in graphene



$$E_{N,\pm} = \pm \sqrt{2\hbar c^2 eB(N + \frac{1}{2} \pm \frac{1}{2})}$$

Level-degeneracy:

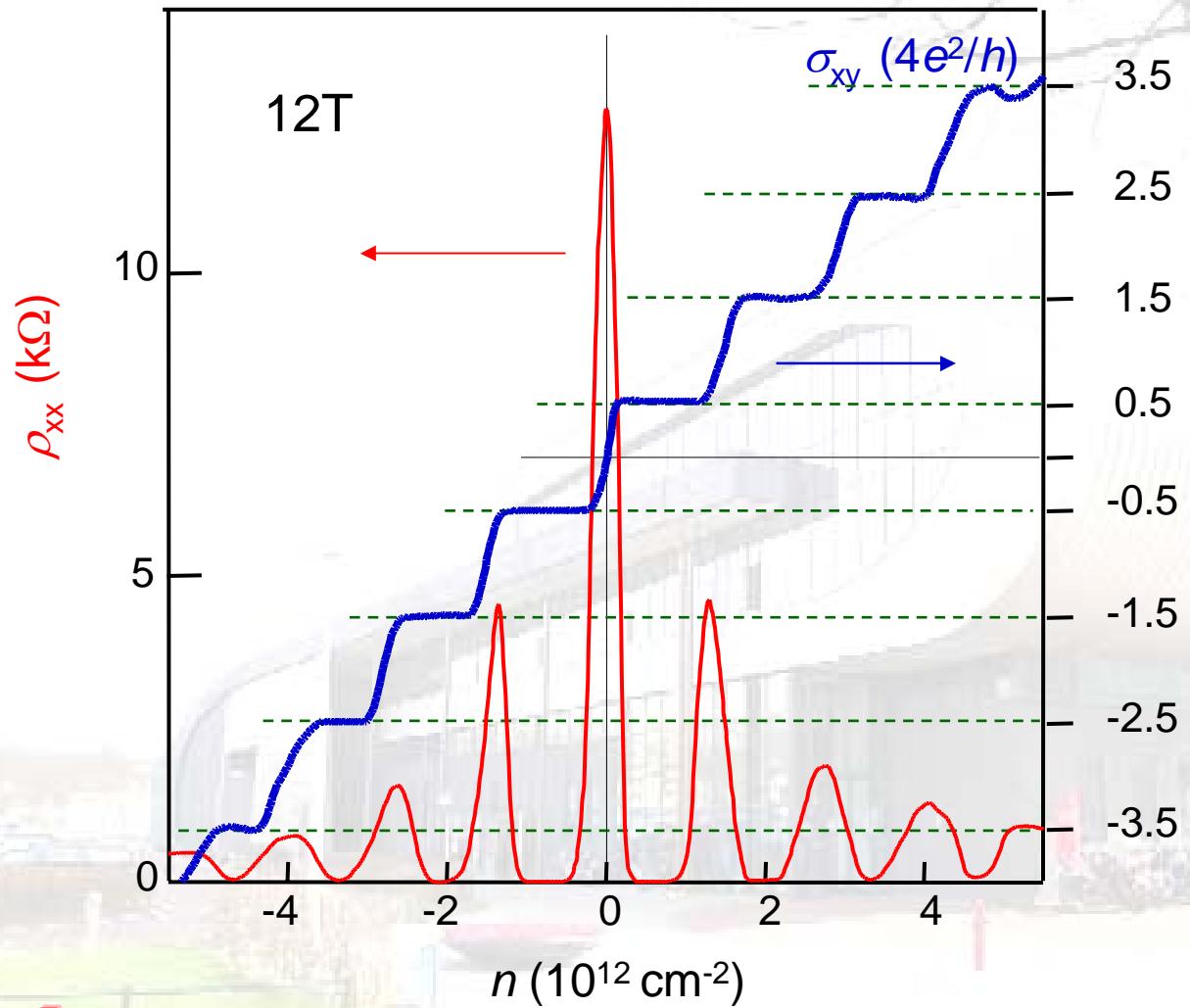
$$n_L = g_v g_s \frac{eB}{h} = 4 \frac{eB}{h}$$

$v = 0$  Landau level **at zero energy**  
*half-filled with electrons and holes*

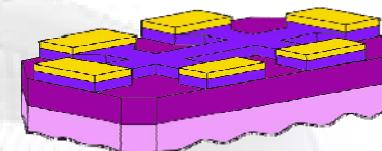
⇒ half-integer QHE



## Half-integer quantum Hall effect



$$\sigma_{xy} = \frac{4 e^2}{h} \left( N + \frac{1}{2} \right)$$

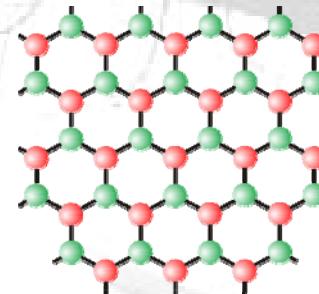
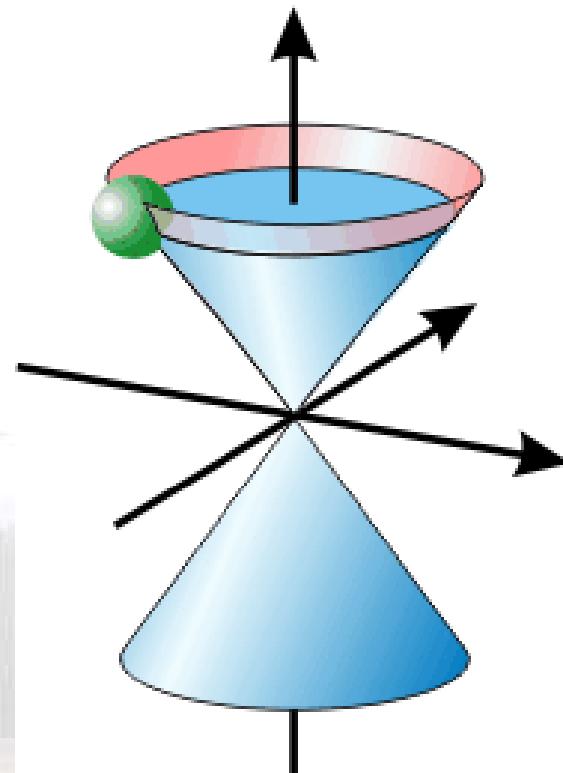


Novoselov *et al*, Nature 2005  
Zhang *et al*, Nature 2005



## Berry's phase

$$\psi = \sin \varphi |+> + \cos \varphi |->$$



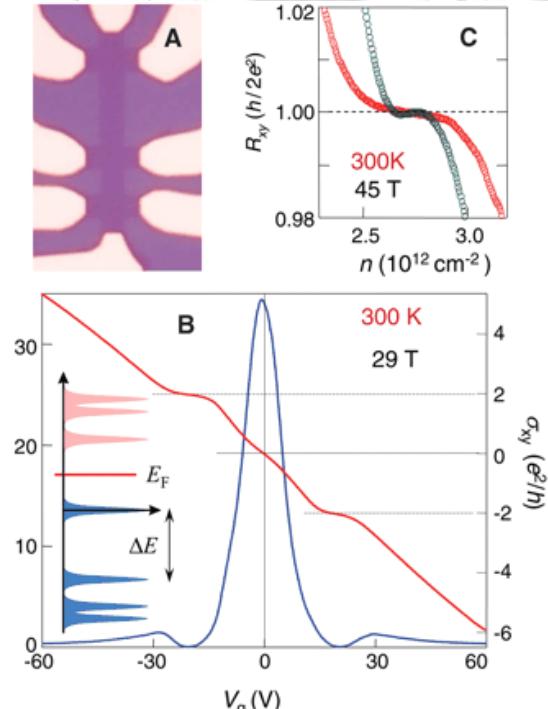
Chiral superposition of  
sublattice wave functions

⇒ Topological  
Berry's phase  $\pi$



## Room Temperature QHE

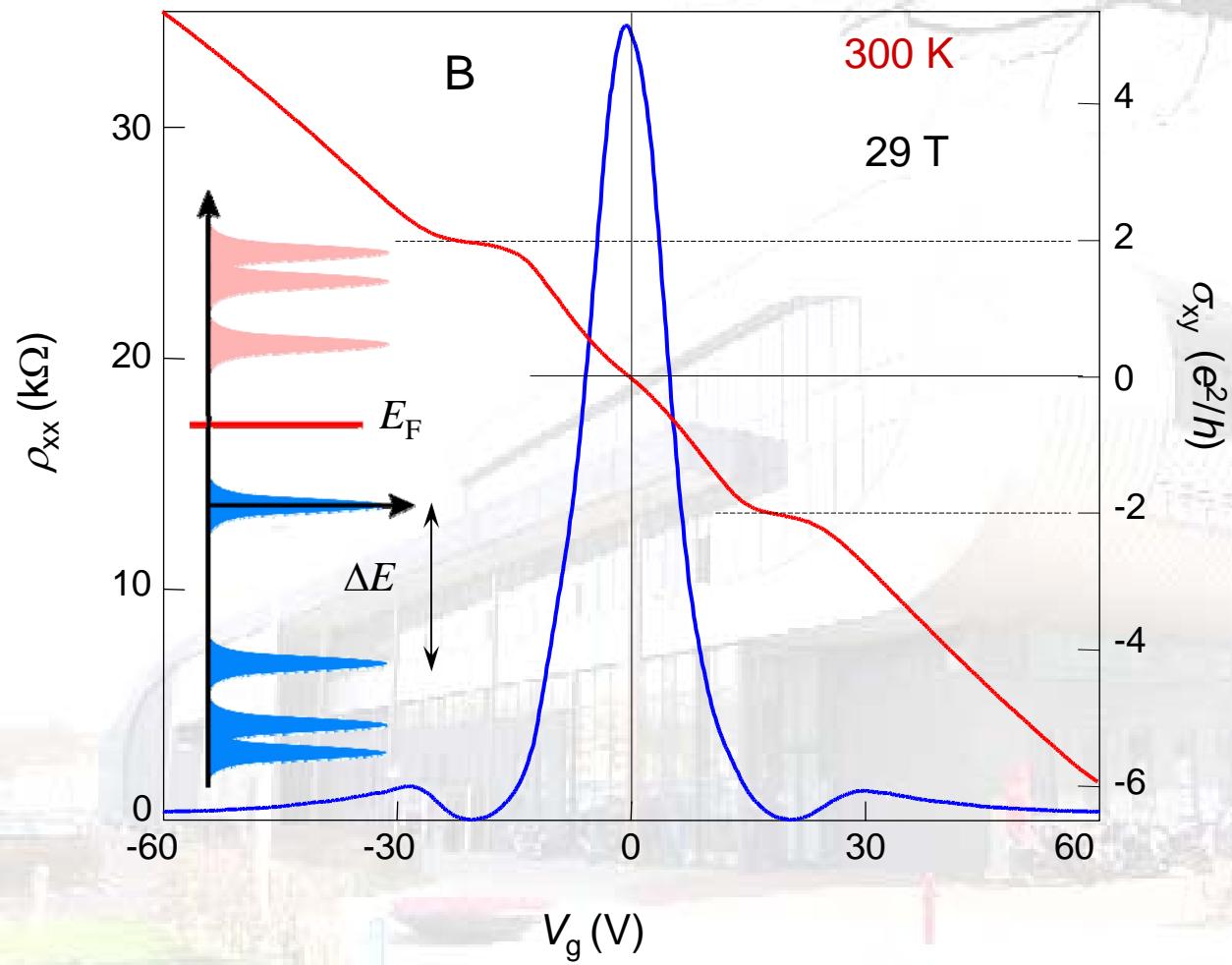
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  - A new type of integer QHE
- Conclusions



K.S. Novoselov *et al.*,  
Science 315, 1379 (2007).



## Room Temperature QHE at 30 T

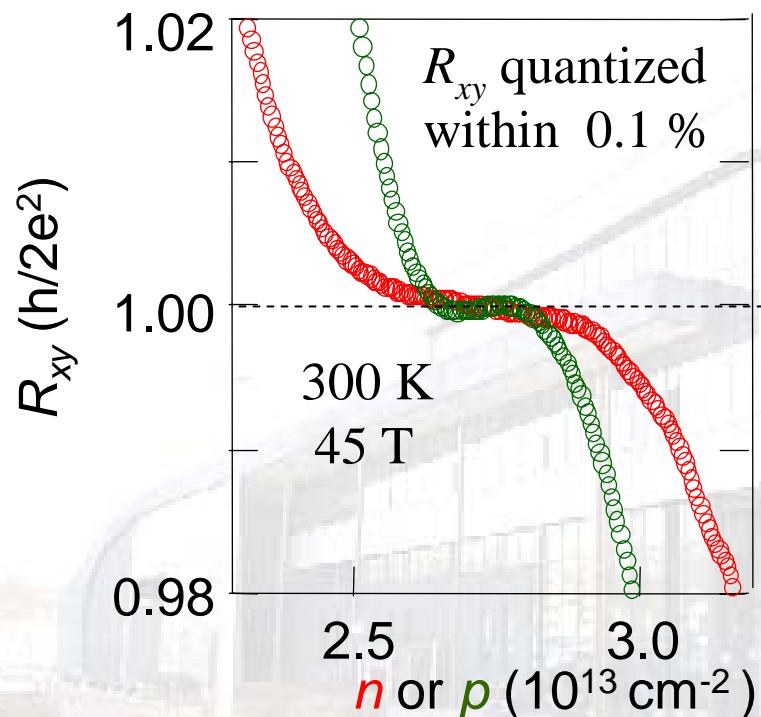


*previously only  
below 30 K !*

Novoselov *et al.*,  
Science 2007



## Room Temperature QHE at 45 T

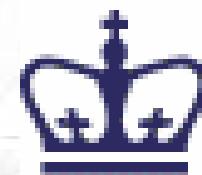


With:

Z. Jiang,\* Y. Zhang, H. L. Stormer,  
G. Boebinger,\* and P. Kim,

*Columbia University, New York, USA*

\**National High Magnetic Field Laboratory,  
Tallahassee, USA*



Novosolov *et al.*, Science 2007





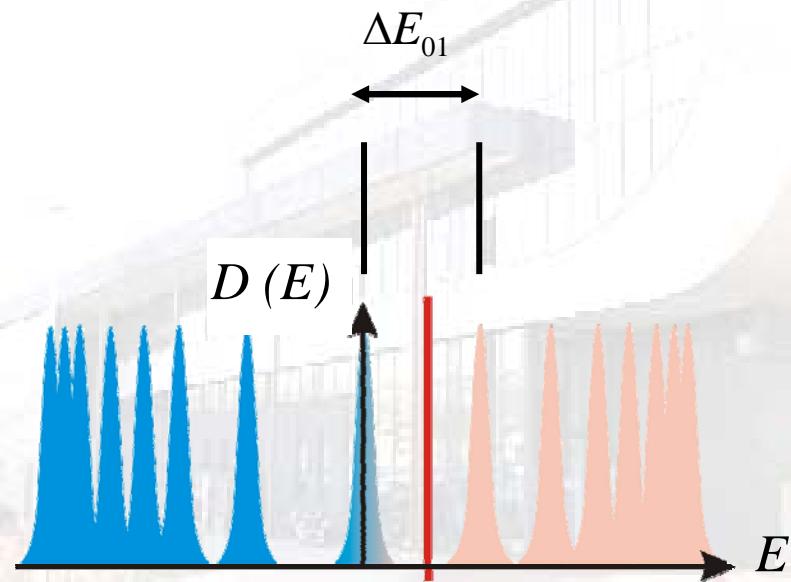
# Gaps in graphene

$$E_{N,\pm} = \text{sgn}(N) \sqrt{2\hbar c^2 e B |N + \frac{1}{2} \pm \frac{1}{2}|}$$

$$\Delta E_{kl} = c_* \sqrt{2e\hbar B} (\sqrt{l} - \sqrt{k})$$

**B = 30 T**

$$\Delta E_{01} \approx 2300 \text{ K}$$





## ... more than just gaps!

- single-layer graphene:

$$E_{01} = \pm \sqrt{2\hbar c^2 e B} \approx 2300 \text{ K} \quad \text{at } 30 \text{ T}$$

we do observe room-temperature QHE !

- conventional semiconductors:

$$\Delta E_{LL} = \frac{\hbar e B}{2m^*} \approx 600 \text{ K} \quad \text{for GaAs}$$

$$\approx 3000 \text{ K} \quad \text{for InSb}$$

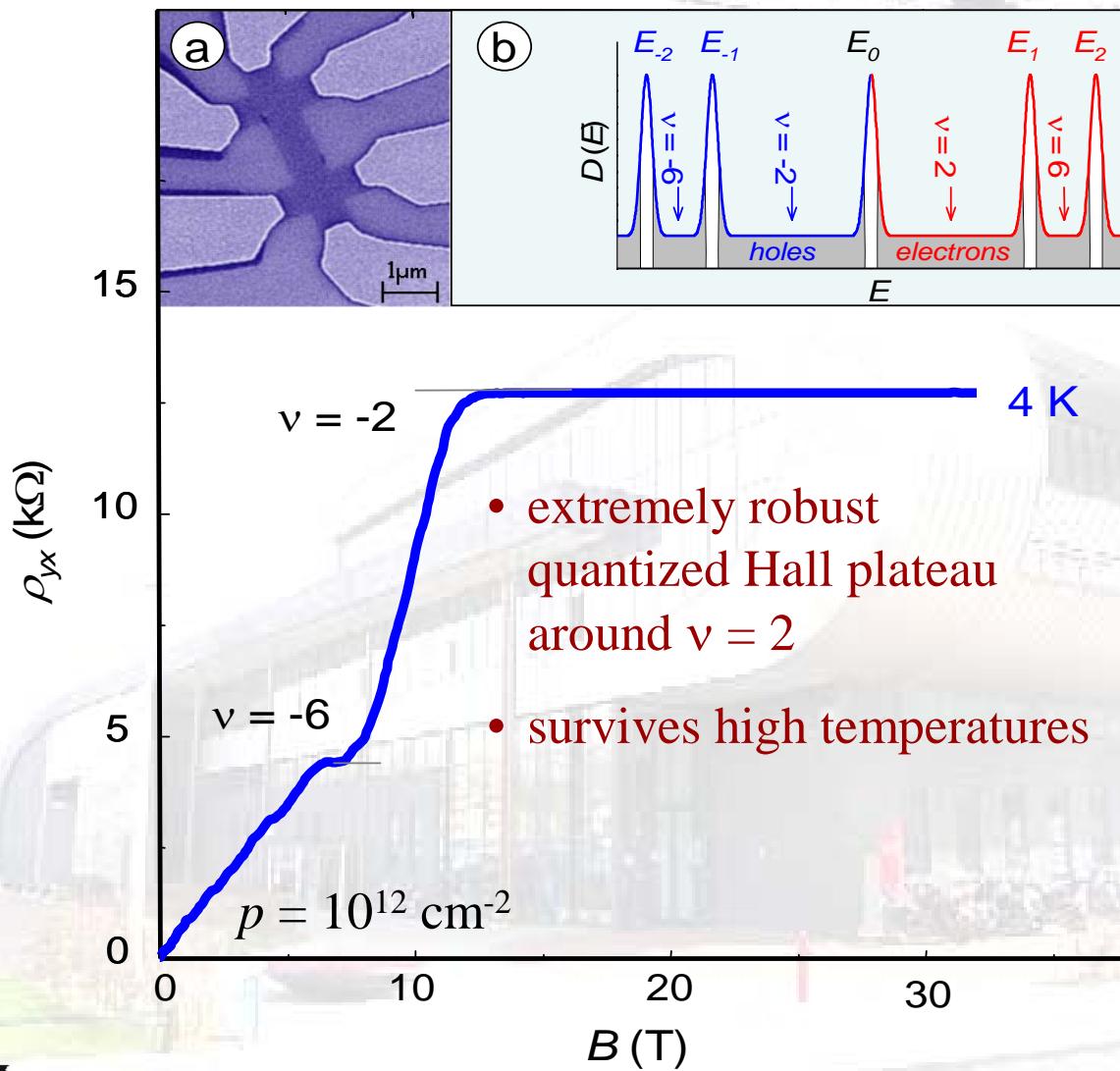
no QHE at 77 K !!

no QHE at 300 K !!

⇒ more than just gaps !!

- *chirality of edge states ?*
- *mesoscopic samples ?*
- *non-local transport ?*
- *localization in graphene !?*
- *peculiar zero-energy level !?*
- .....???

# Robustness of the QHE





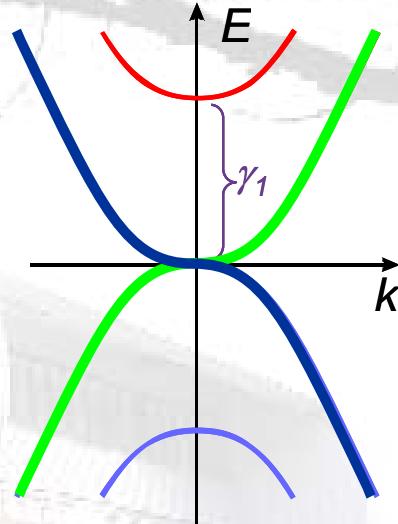
# Bilayer graphene

- Introduction
  - Graphene devices
  - Quantum Hall effect
- Single-layer graphene
  - Bandstructure of graphene
  - Half-integer QHE of chiral Dirac fermions
  - Room-temperature QHE & activation gaps

## ➤ Bilayer graphene

### ➤ Bandstructure of massive chiral Dirac fermions

- A new type of integer QHE
- Conclusions

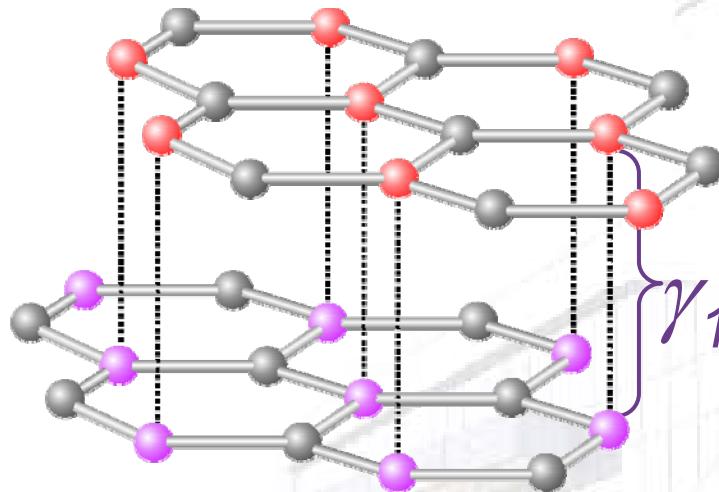


with: Ed McCann and Vladimir Fal'ko

*Department of Physics, Lancaster University, UK*

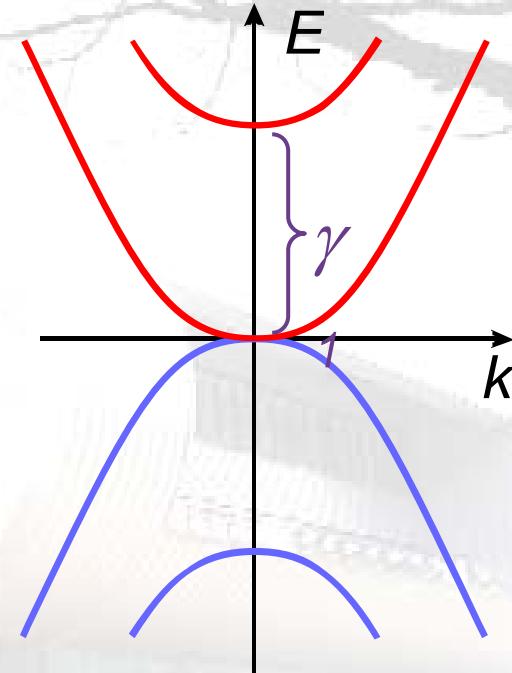


# Bandstructure in bilayers



A-B stacking

$\gamma_1$ : inter-layer coupling  
 $c$ : particle velocity



$$E(p) = \pm \left( \frac{1}{2} \gamma_1 + \sqrt{\frac{1}{4} \gamma_1^2 + c^2 p^2} \right)$$

McCann, Fal'ko, PRL **96**, 086805 (2006)



# QED in a pencil trace

$$E(p) = \pm \frac{1}{2} \gamma_1 \pm \sqrt{\frac{1}{4} \gamma_1^2 + v_F^2 p^2}$$

$$E(p) = \pm \left( m_0 c^2 + \sqrt{m_0^2 c^4 + p^2 c^2} \right)$$

inter-layer

intra-layer

massive  
relativistic  
chiral

} Dirac fermions



***High energy physics  
in a pencil trace***

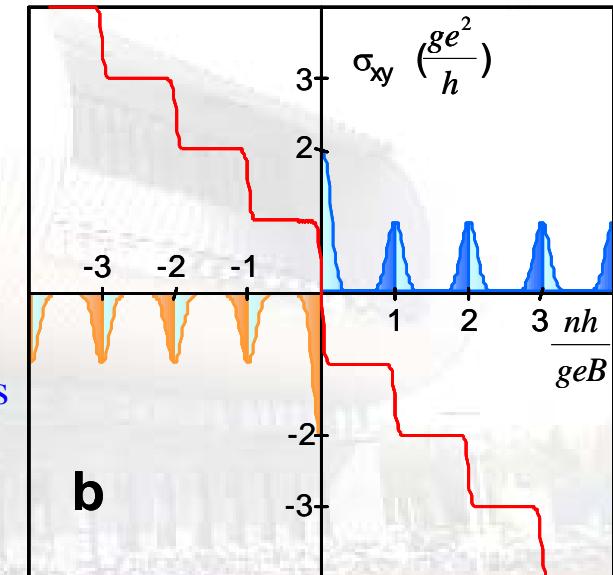




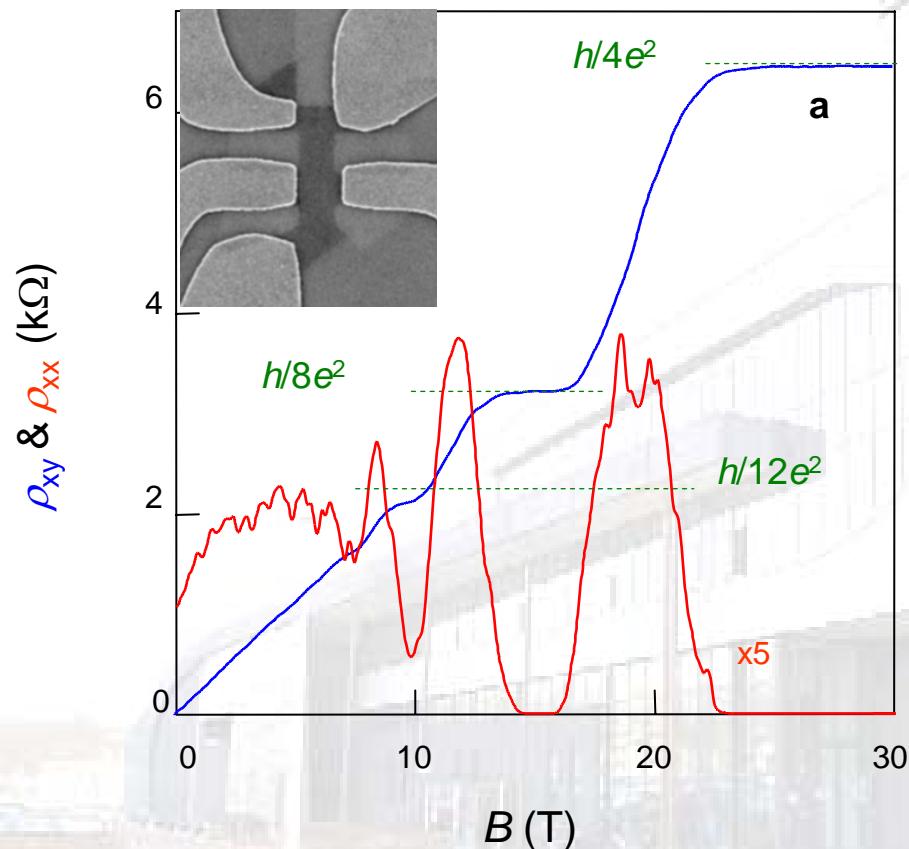
# A new type of integer QHE

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  - Room-temperature QHE & activation gaps
- **Bilayer graphene**
  - Bandstructure of massive chiral Dirac fermions
  - **A new type of integer QHE**
- Conclusions

**K.S. Novoselov *et al.*,**  
**Nature Physics 2, 177 (2006).**



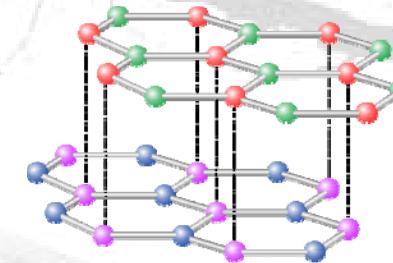
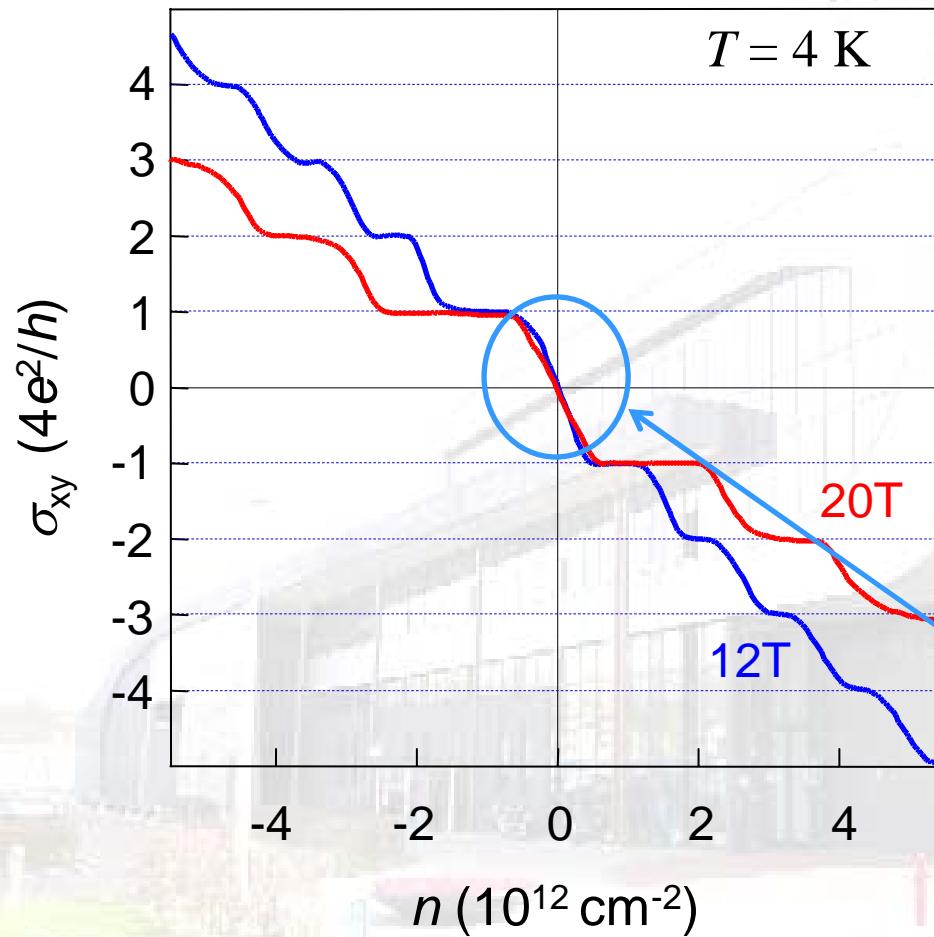
## QHE in bilayer graphene



- ⇒ re-establishment of integer QHE  
(four-fold degenerate Landau levels)
- ⇒ towards a conventional QHE  
with a parabolic dispersion ?

↑ Novoselov et al, Nature Physics 2006

## Unconventional integer QHE



- quantum Hall plateaus  
(of four-fold degenerate Landau levels)  
at integer filling factors  
 $N = \pm 1, \pm 2, \pm 3, \dots$

**$N=0$  plateau missing !**



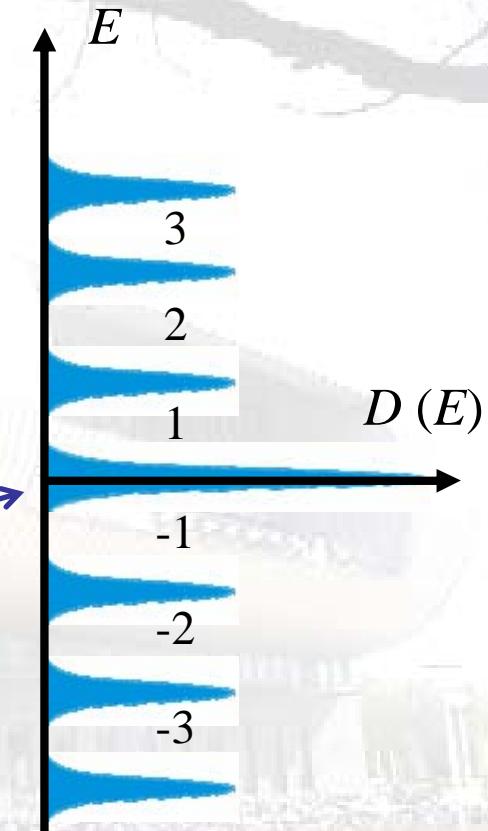
## Landau levels in a bilayer

$$E_{LL} = \pm \hbar \omega_c \sqrt{N(N-1)}$$

$$= \begin{cases} 0 & \text{for } N = 0 \\ \pm \hbar \omega_c (N + \frac{1}{2}) & \text{for } N \rightarrow \infty \end{cases}$$

**8-fold degenerate** Landau level  
at zero energy shared equally  
between electrons and holes

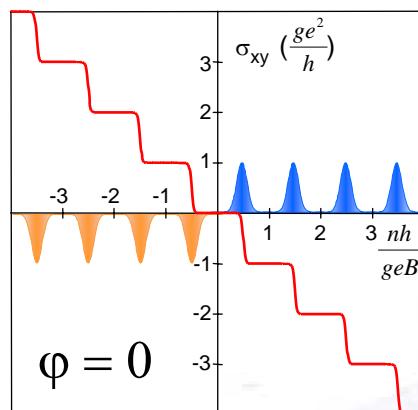
topological (Berry) phase  $2\pi$





## Three types of integer QHE

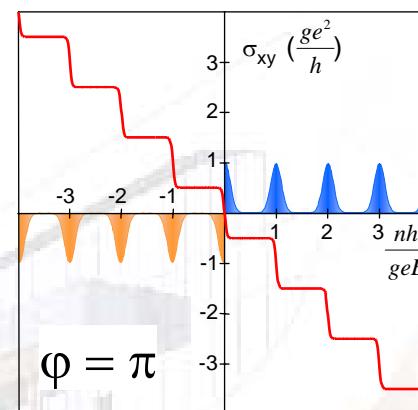
conventional  
integer QHE  
(semiconductors)



$$E_{LL} = \pm \hbar \omega_c (n + \frac{1}{2})$$

von Klitzing *et al.*  
PRL **45**, 494 (1980)

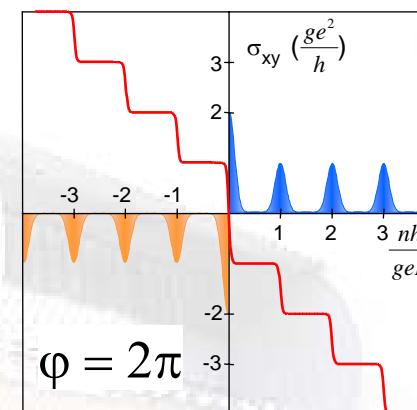
relativistic  
half-integer QHE  
(single-layer graphene)



$$E_{LL} = \pm \sqrt{2\hbar c^2 e B (n + \frac{1}{2} \pm \frac{1}{2})}$$

Novoselov *et al.*,  
Nature **438**, 197 (2005)  
Zhang *et al.*,  
Nature **438**, 201 (2005)

unconventional  
integer QHE  
(bilayer graphene)



$$E_{LL} = \pm \hbar \omega_c \sqrt{n(n-1)}$$

Novoselov *et al.*,  
Nature Physics **2**, 177 (2006).  
McCann, Fal'ko,  
PRL **96**, 086805 (2006)

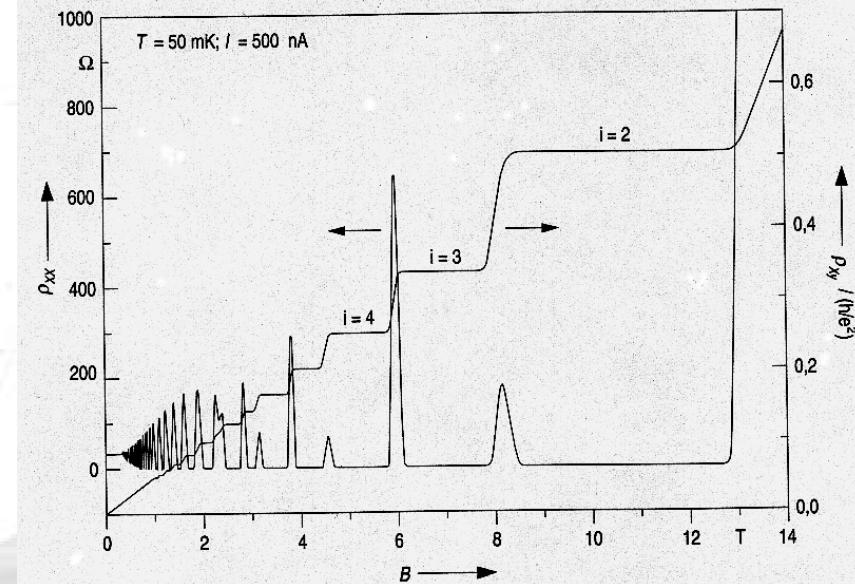
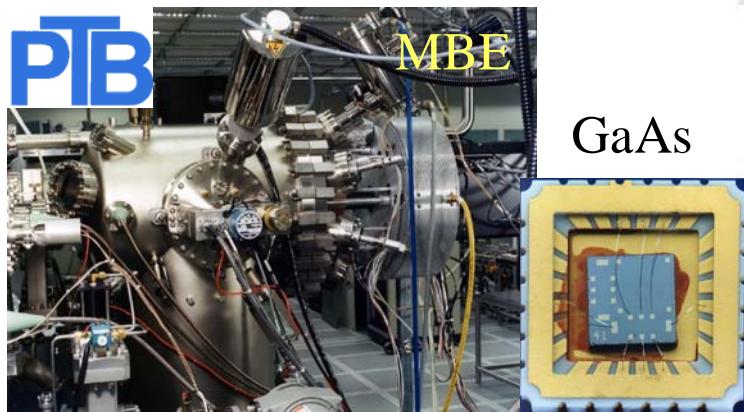




QHE in Graphene

## QHE today

PB



$$R_K = 25812.807 \Omega$$



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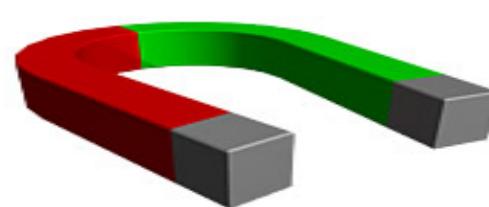
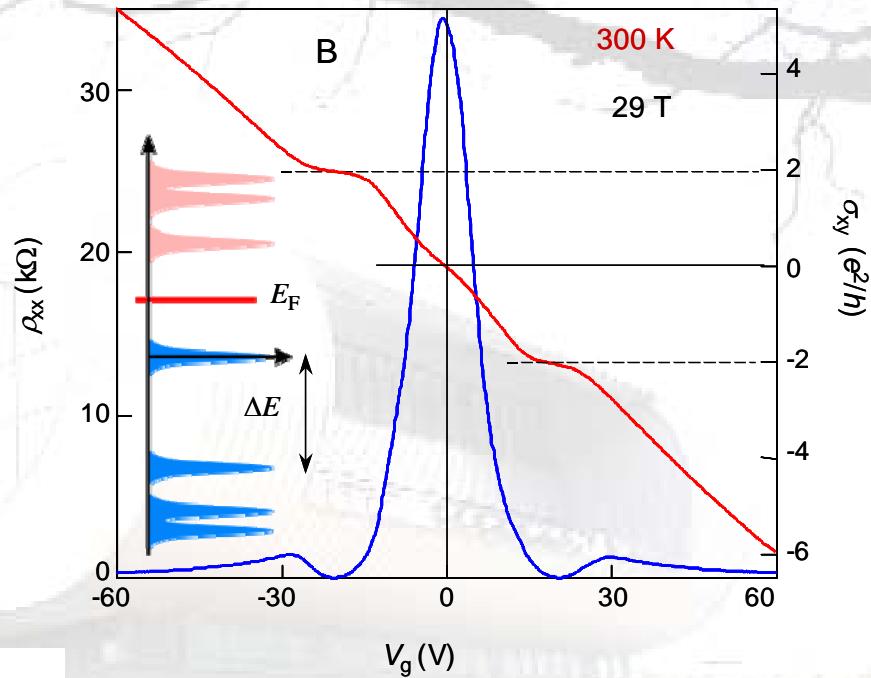


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QHE in Graphene

# Future of QHE



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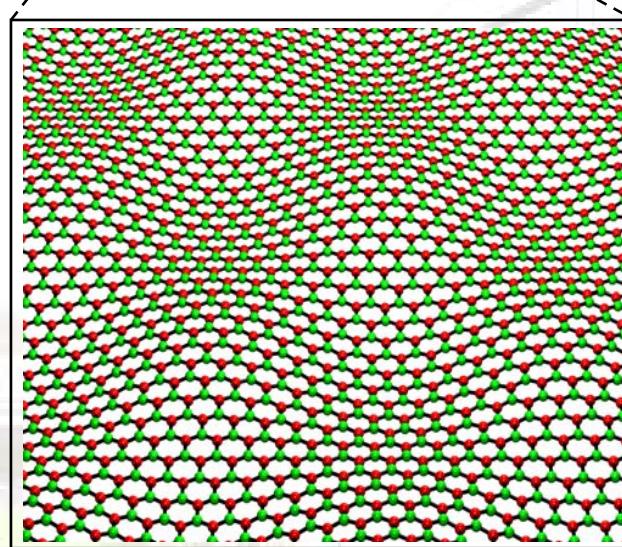
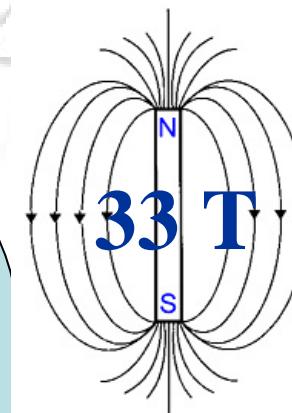
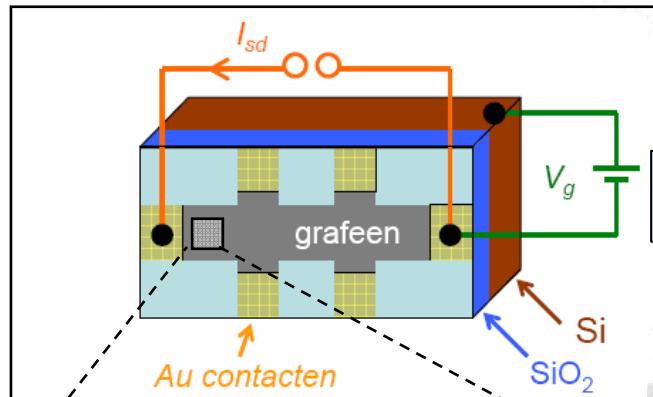
# Conclusions

## Conclusions

- **Graphene:** a real two-dimensional crystal !
- ***single-layer graphene:***
  - relativistic electrons with a half-integer QHE
  - room-temperature QHE
  - towards a room-temperature quantum metrology ?
- ***bilayer graphene:*** a third type of QHE with chiral, massive relativistic Dirac fermions
- **after 25 years QHE** is still alive
- **high magnetic fields** can help a lot



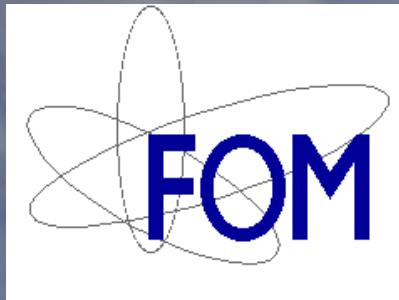
# Commercial





QHE in Graphene

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## News from the QHE in Graphene



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Condensed Matter Science / High Field Magnet Laboratory  
Institute for Molecules and Materials, Radboud University Nijmegen

Review: A.K. Geim and K.S. Novoselov, *The rise of graphene*,  
Nature Materials **6**, 183-191 (March 2007).

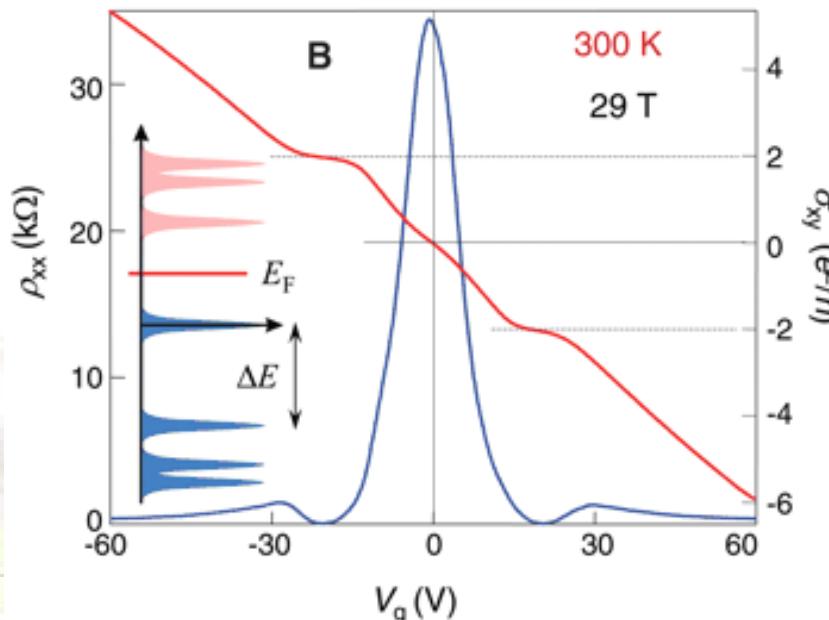
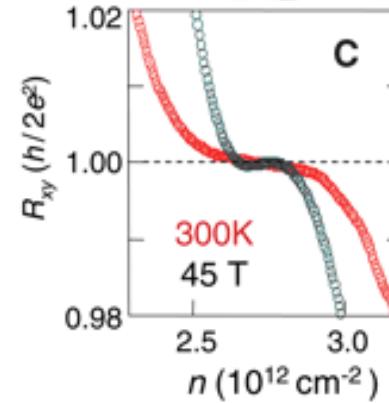


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## Room-temperature QHE in graphene



K.S. Novoselov *et al.*,  
Science 315, 1379 (2007).

