

Tunneling in suspended carbon nanotubes assisted by longitudinal phonons

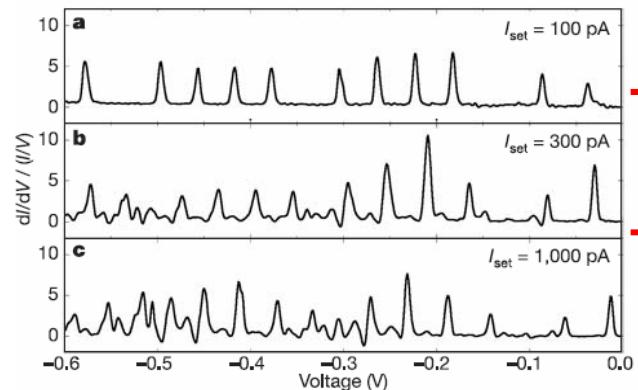
S. Sapmaz, P. Jarrillo-Herrero, Ya. M. Blanter,
C. Dekker, and H. S. J. Van der Zant

Cond-mat/0508270

Current-voltage characteristics of suspended single-wall carbon nanotube quantum dots show a series of steps equally spaced in voltage. The energy scale of this harmonic, low-energy excitation spectrum is consistent with that of the longitudinal low- k phonon mode (stretching mode) in the nanotube. Agreement is found with a Franck-Condon-based model in which the phonon-assisted tunneling process is modeled as a coupling of electronic levels to underdamped quantum harmonic oscillators. Comparison with this model indicates a rather strong electron-phonon coupling factor of order unity.

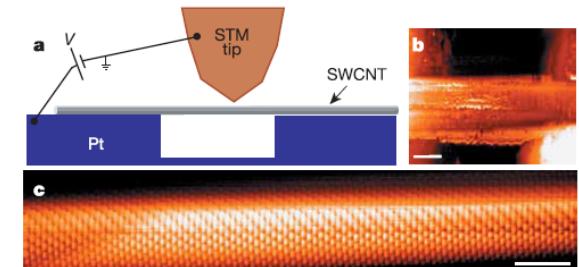
Vibrational effects at the molecular scale

Electrical generation and absorption of phonons in CNTs
B. J. LeRoy *et al.*, Nature 432, 371 (2004)



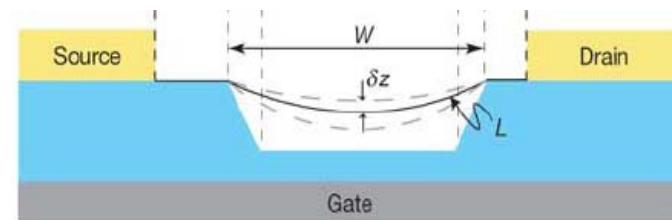
→ Coulomb peaks

→ Coulomb+phonon satellites



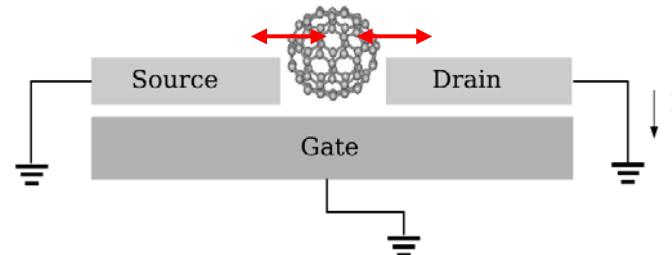
A tunable carbon nanotube electromechanical oscillator

V. Sazonova *et al.*, Nature 431, 284 (2004)



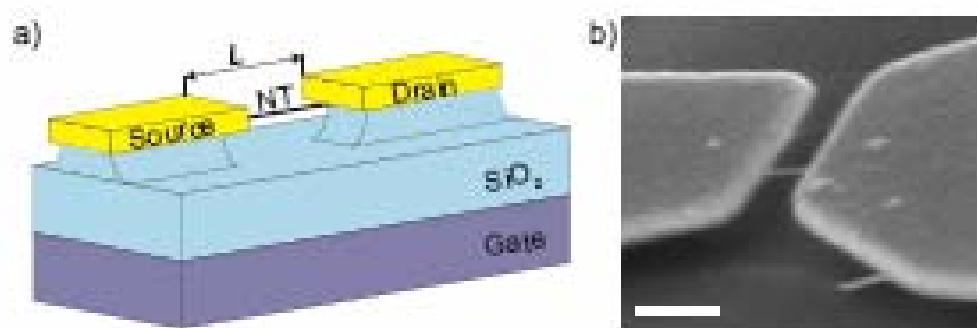
C₆₀ shuttle

Park *et al.*, Nature 407, 57 (2000)



Device fabrication

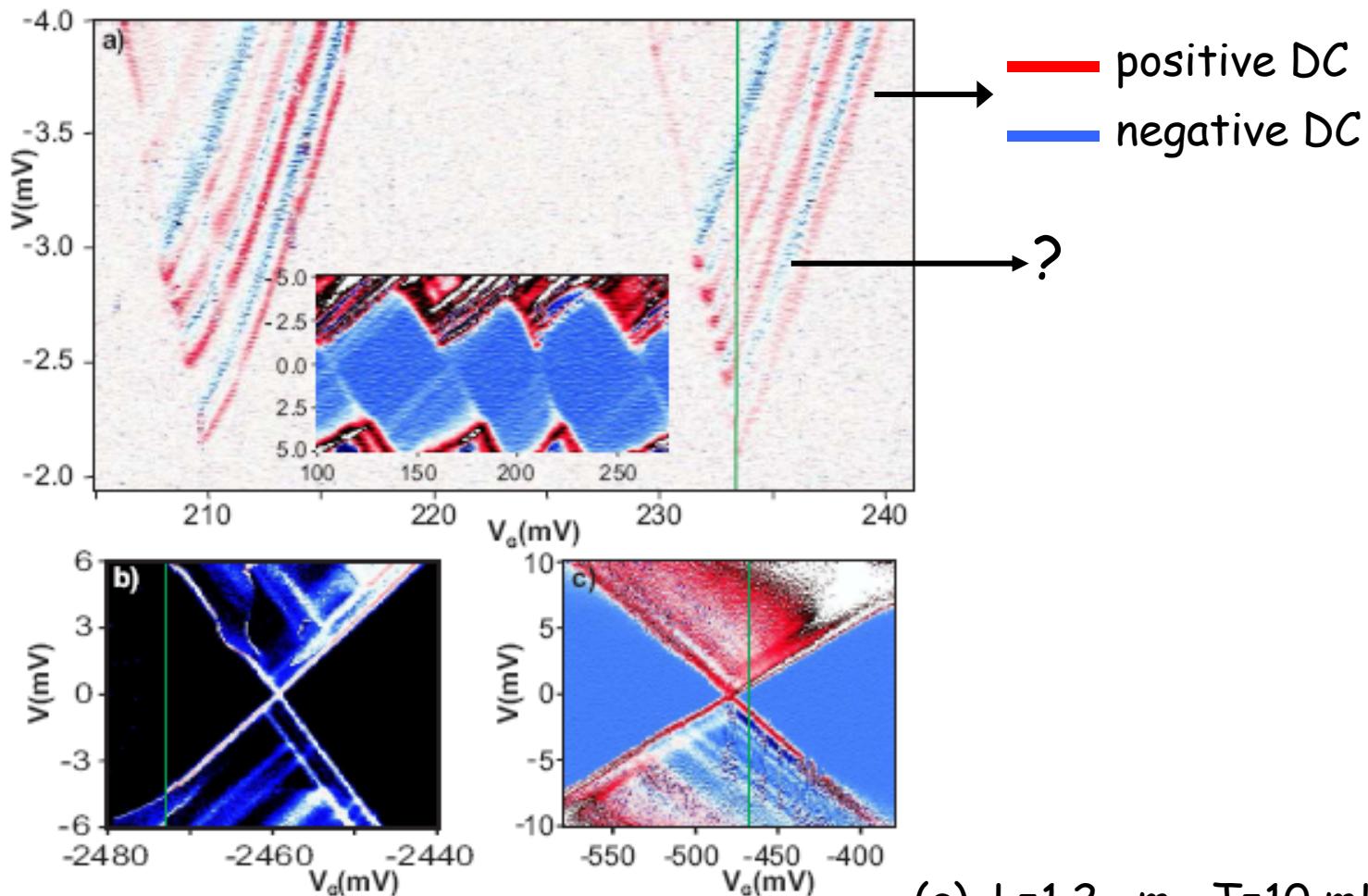
Cond-mat/0508270



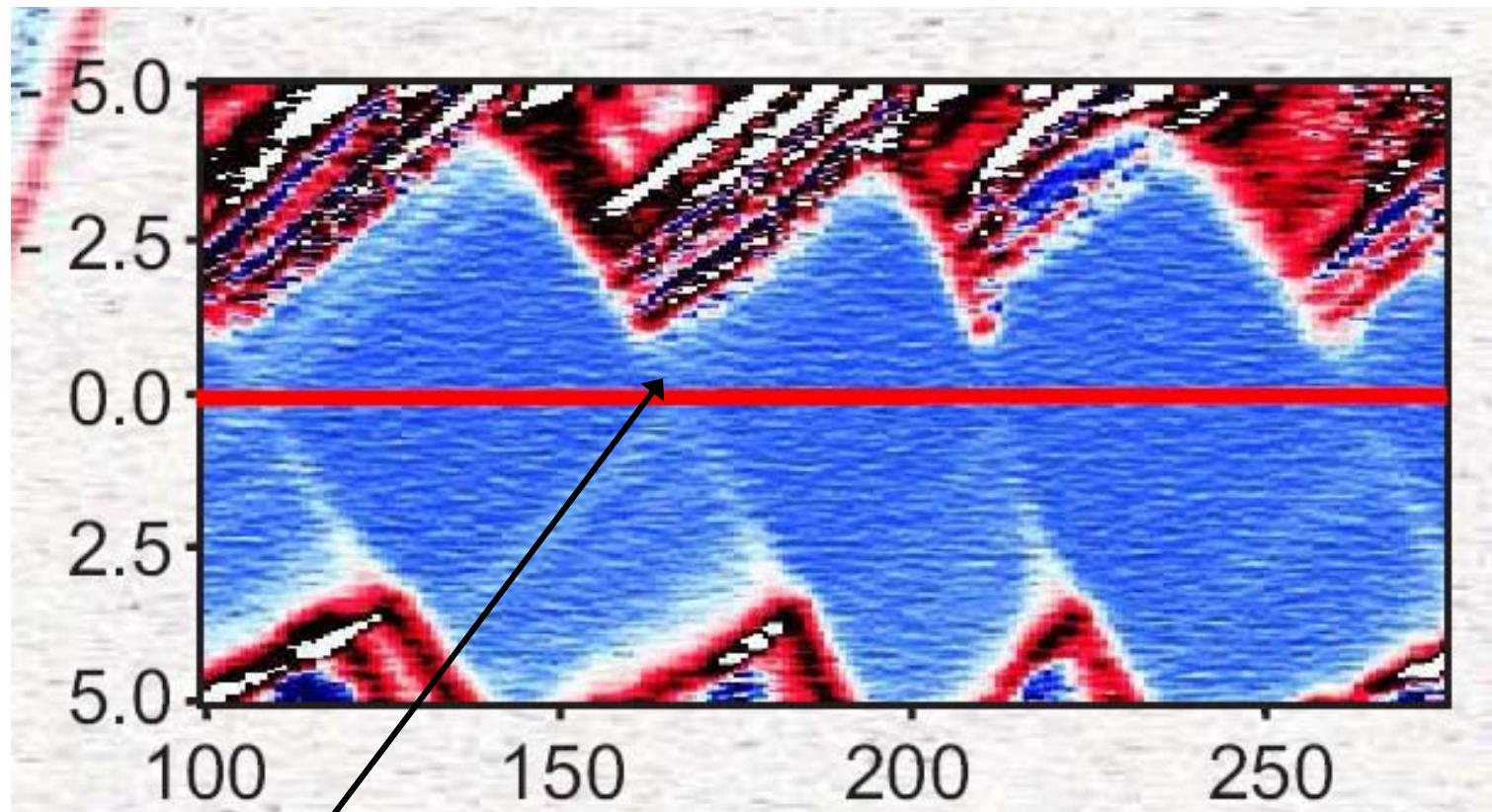
SEM micropgraph

- Individual CNTs produced by laser ablation and CVD
- CNTs on Si/SiO₂ substrate via AFM
- Electrodes via e-beam lithography + thermal evaporation of Cr/Au
- Remove underlying SiO₂ via buffered HF

Stability diagrams

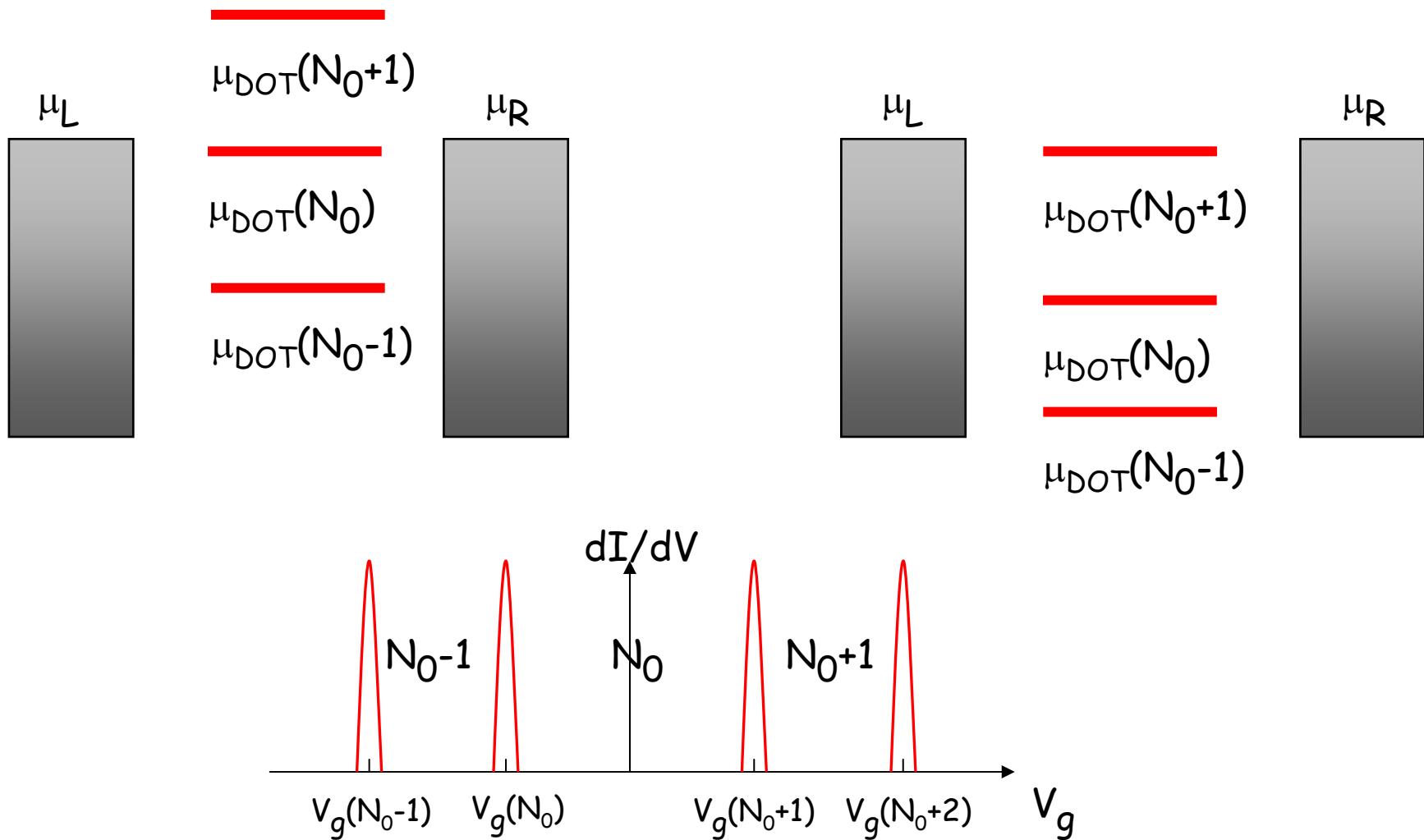


- (a) $L=1.2 \mu\text{m}$, $T=10 \text{ mK}$ (inset 300 mK)
- (b) $L=420 \text{ nm}$, $T=300 \text{ mK}$
- (c) $L=140 \text{ nm}$, $T=300 \text{ mK}$
 $d \sim 1-1.4 \text{ nm}$



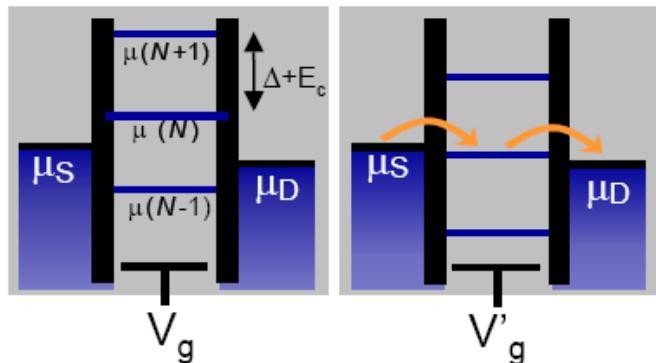
current suppression cannot be counteracted
by changing the gate voltage
Franck-Condon blockade?

Coulomb blockade in quantum dots

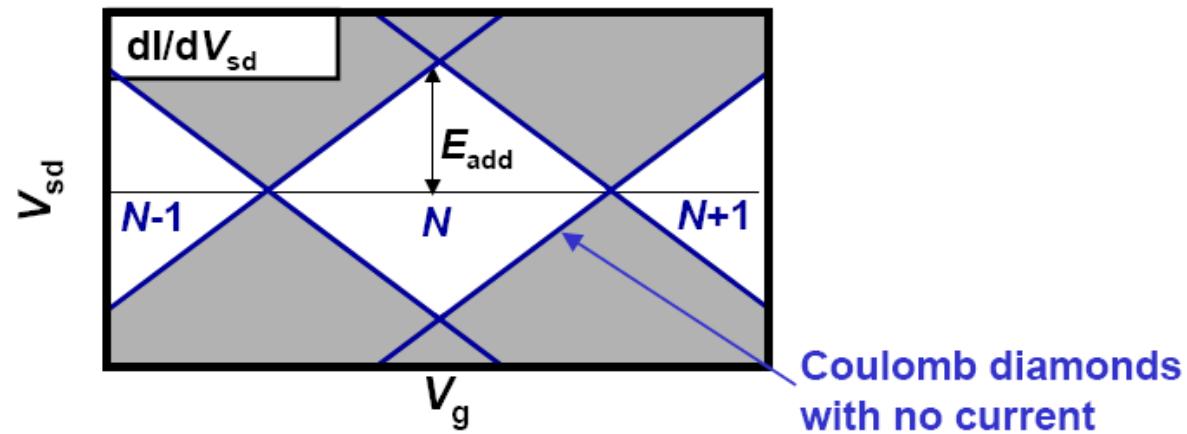


stability diagram and Coulomb diamonds

linear regime

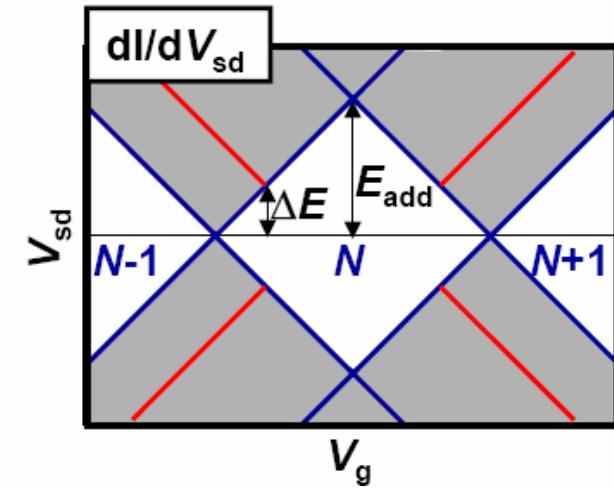
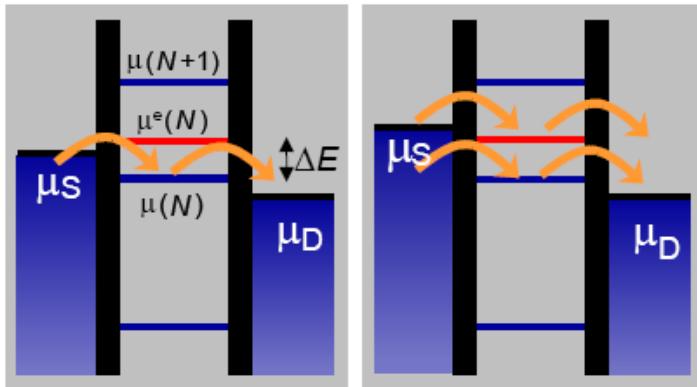


- two 'knobs': V_{SD} and V_{gate}
- state contributes to transport if $\mu_s > \mu(N) > \mu_d$

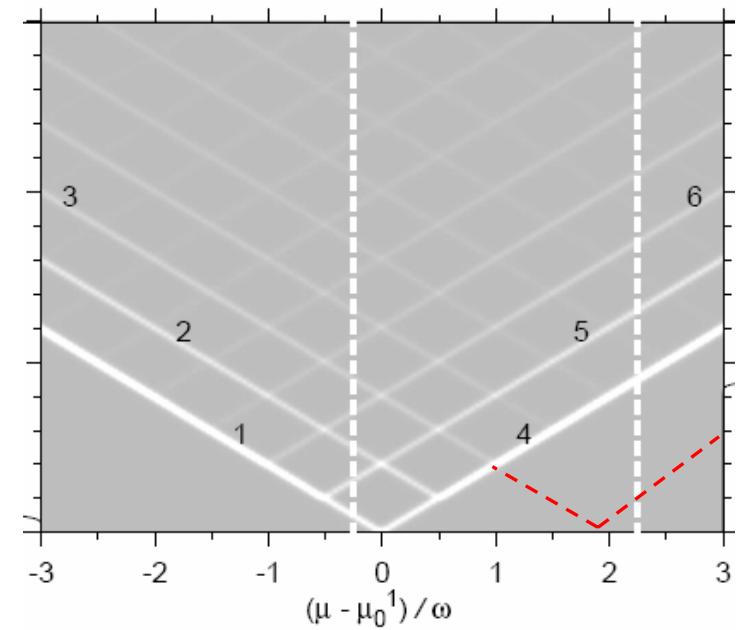


from: Herre van der Zant's talk

level spectroscopy

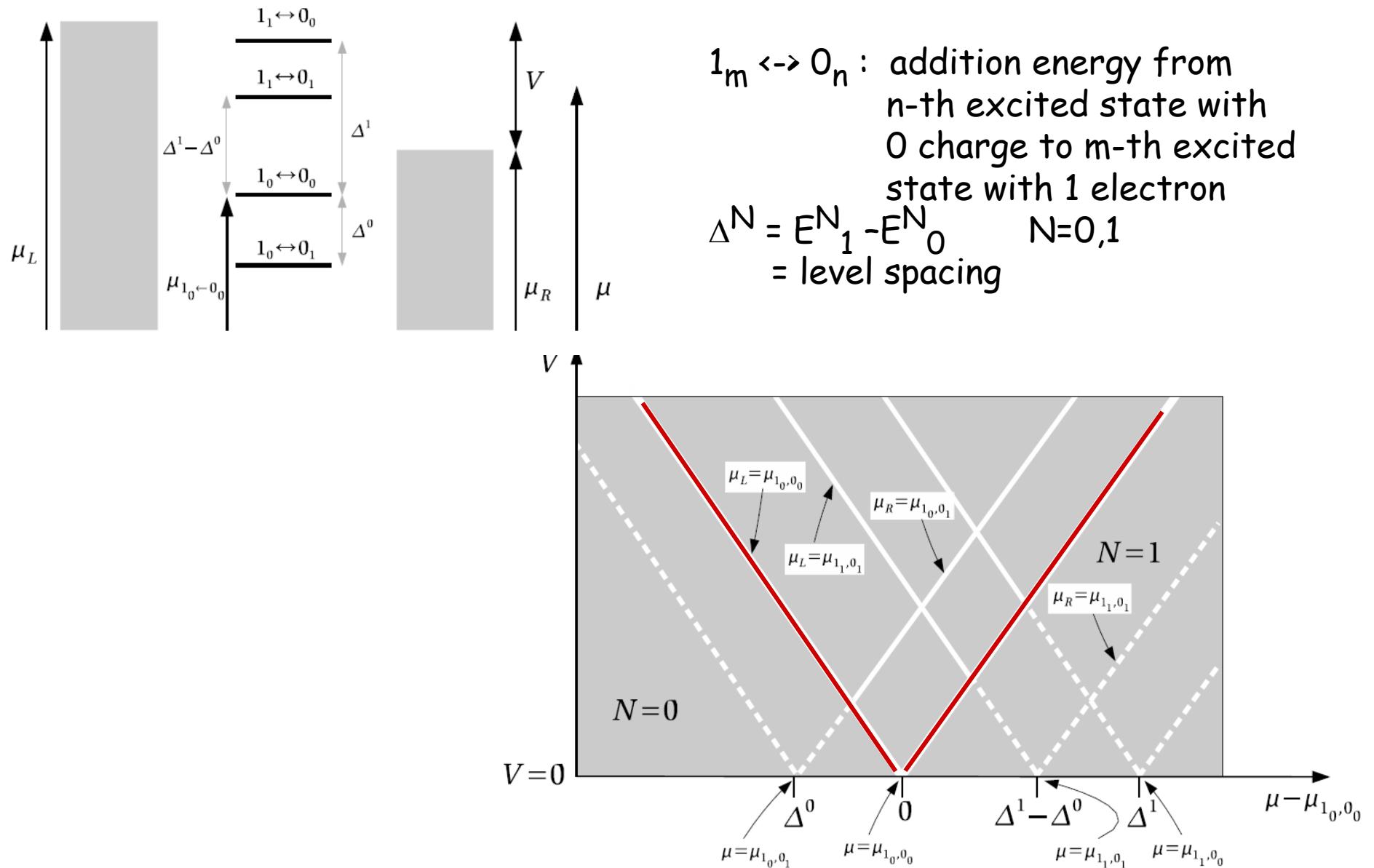


- **excited states result in extra lines in stability diagram (red)**
- **excited states can be electronic, vibrational, or spin related**

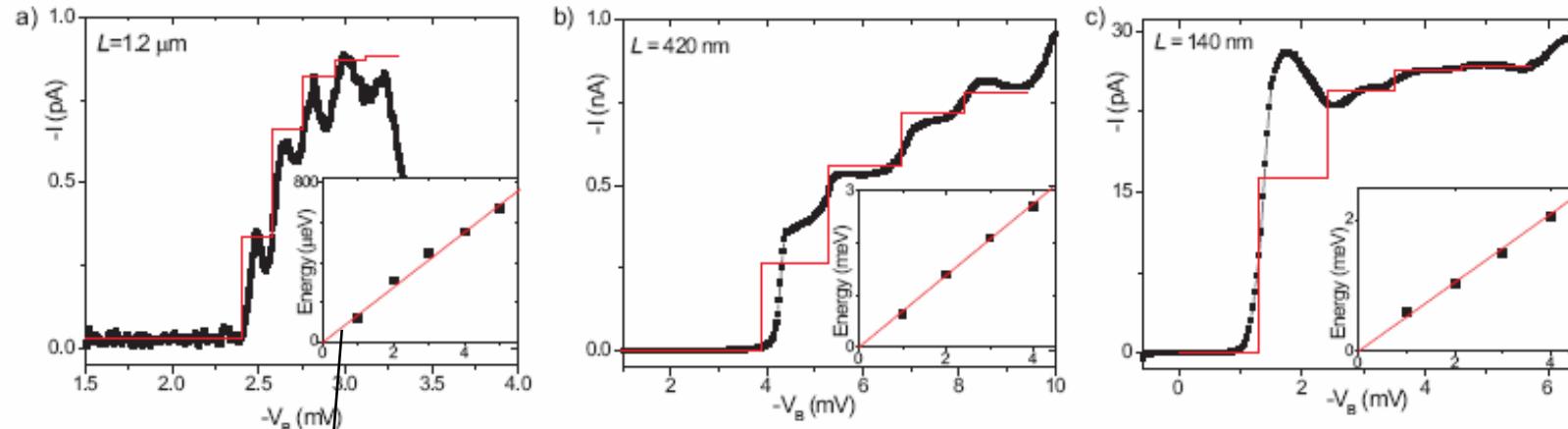


Side bands related to excitations 2 excited states for N=0 and 1

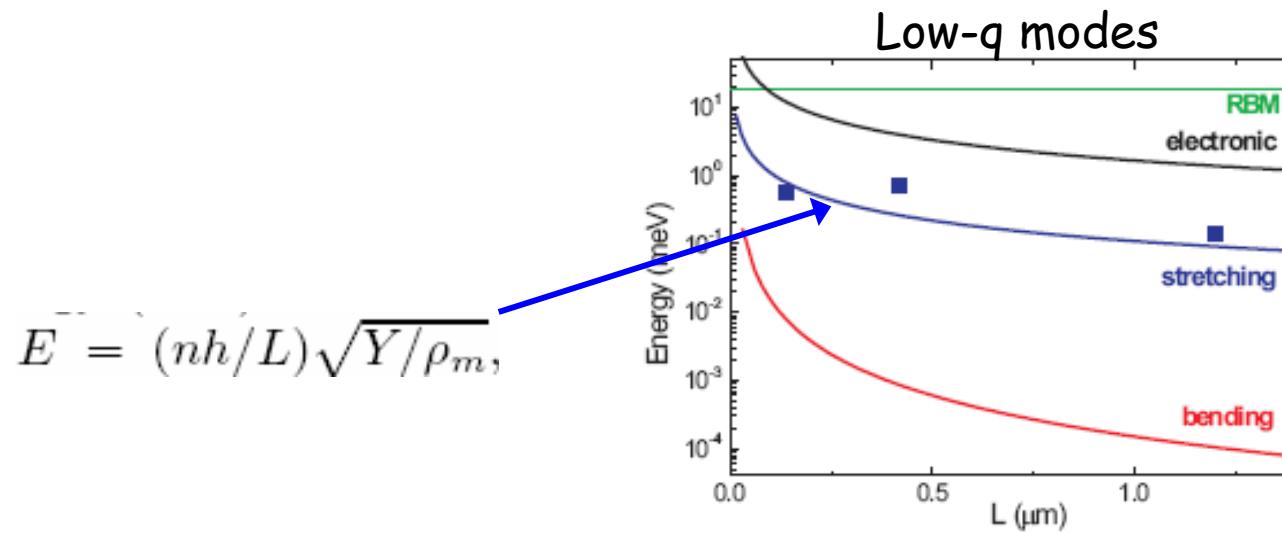
K. Nowack, Diploma (2004)



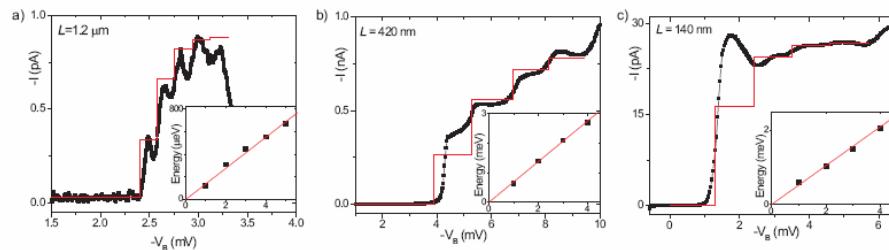
Main point: Energy spacing of excitation lines is a **harmonic spectrum** with fundamental energy $\hbar\Omega$ betw. $140 \mu\text{eV}$ and $530 \mu\text{eV}$ \ll mean electronic level spacing $\Delta \sim \hbar v_F / 2L$



$E \sim \hbar\Omega n$ → Vibrational mode coupled to the tunneling charge
 $n = \text{number of excited vibrational quanta}$
Which mode ???



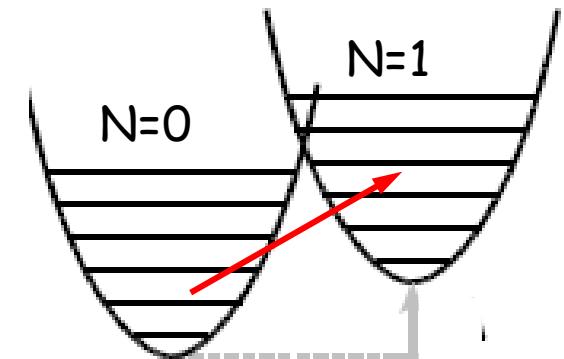
$$E = (nh/L)\sqrt{Y/\rho_m},$$



Franck-Condon (FC) model:

Electronic states coupled to a local vibrational mode

$$\begin{aligned} H &= \varepsilon_d d^\dagger d + H_{Coul} + (\Omega/2)(P^2 + Q^2) + g Q d^\dagger d \\ &= \varepsilon_d d^\dagger d + H_{Coul} + (\Omega/2)(P^2 + [Q + \sqrt{2}(g/\Omega)d^\dagger d]^2) - (...) \end{aligned}$$



transition rates \sim FC factors
nuclear geometry change upon
charging

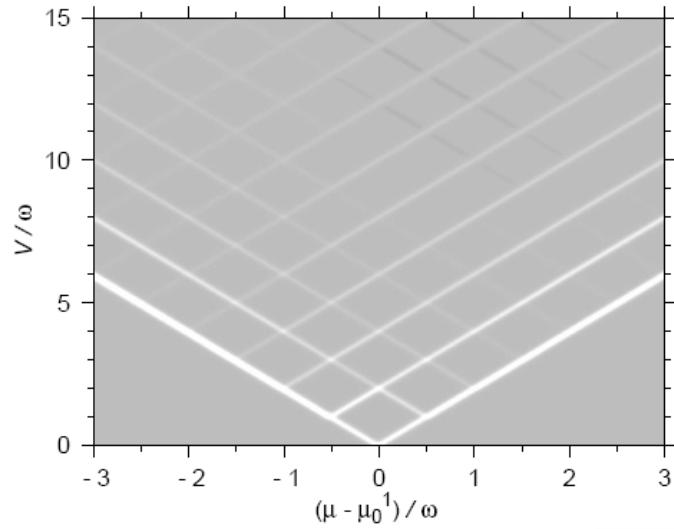
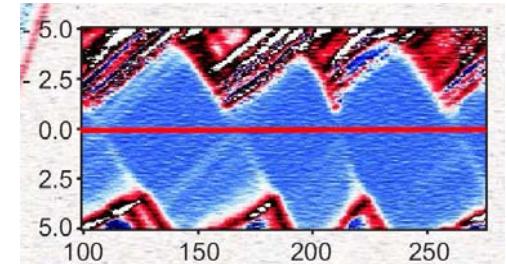
$$F_{q\tilde{q}}^i = \left| \int \Psi_q(Q)^* \Psi_{\tilde{q}}(Q + \sqrt{2} \mathbf{g}) dQ \right|^2$$

Current steps related to \mathbf{g}

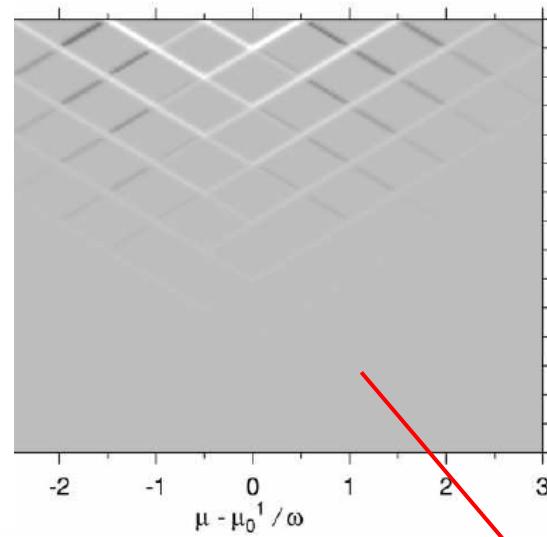
Estimates: $\mathbf{g} \sim 0.95, 1.1$ and 0.5

Franck-Condon blockade for strong electron-vibron coupling

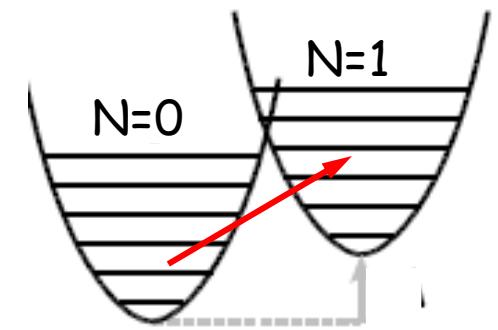
Nowack/Wegewijs, cond-mat/0506552
Koch/von Oppen, PRL 94, 206804 (2004)



$g \ll 1$

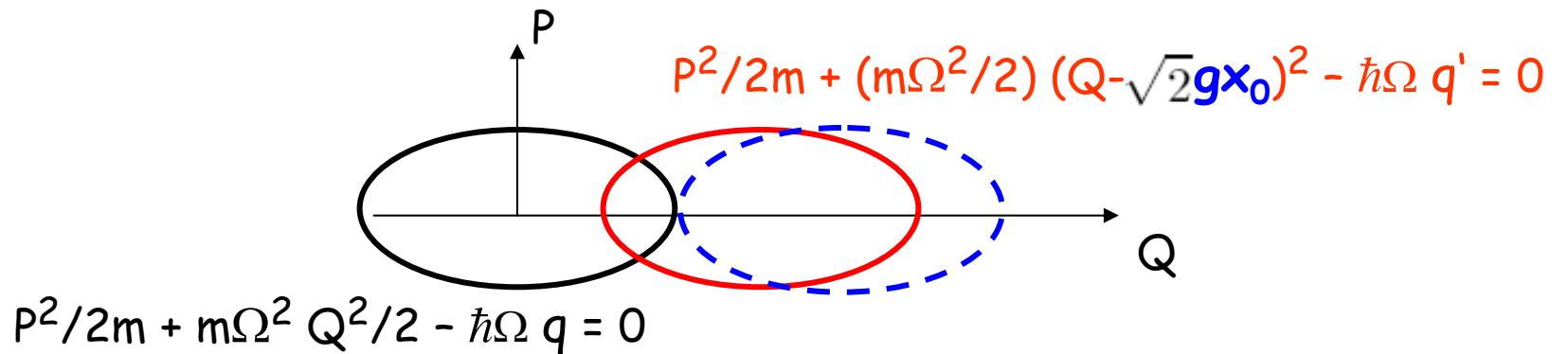


$g \gg 1$



Exponential suppression of
the current at low voltages

Quasi-classical perspective

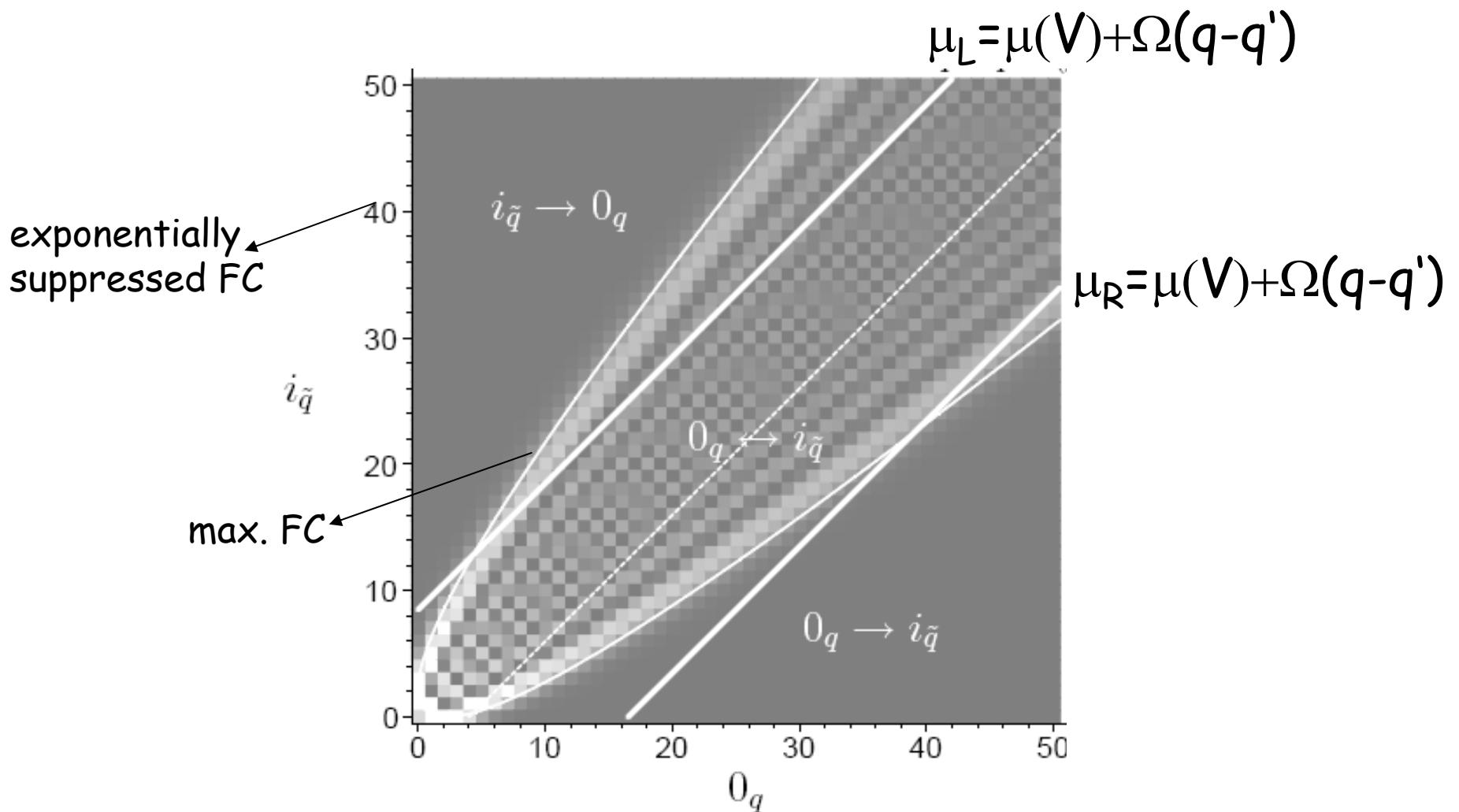


Transitions are ONLY allowed if their **classical** counterparts are also allowed
→ if phase space trajectories of the oscillators do intersect or are tangential !

Condition ($P^2 \geq 0$!)

$$q+q' \geq (q-q')^2/2g^2 + g^2/2 \quad \text{Franck-Condon parabola}$$

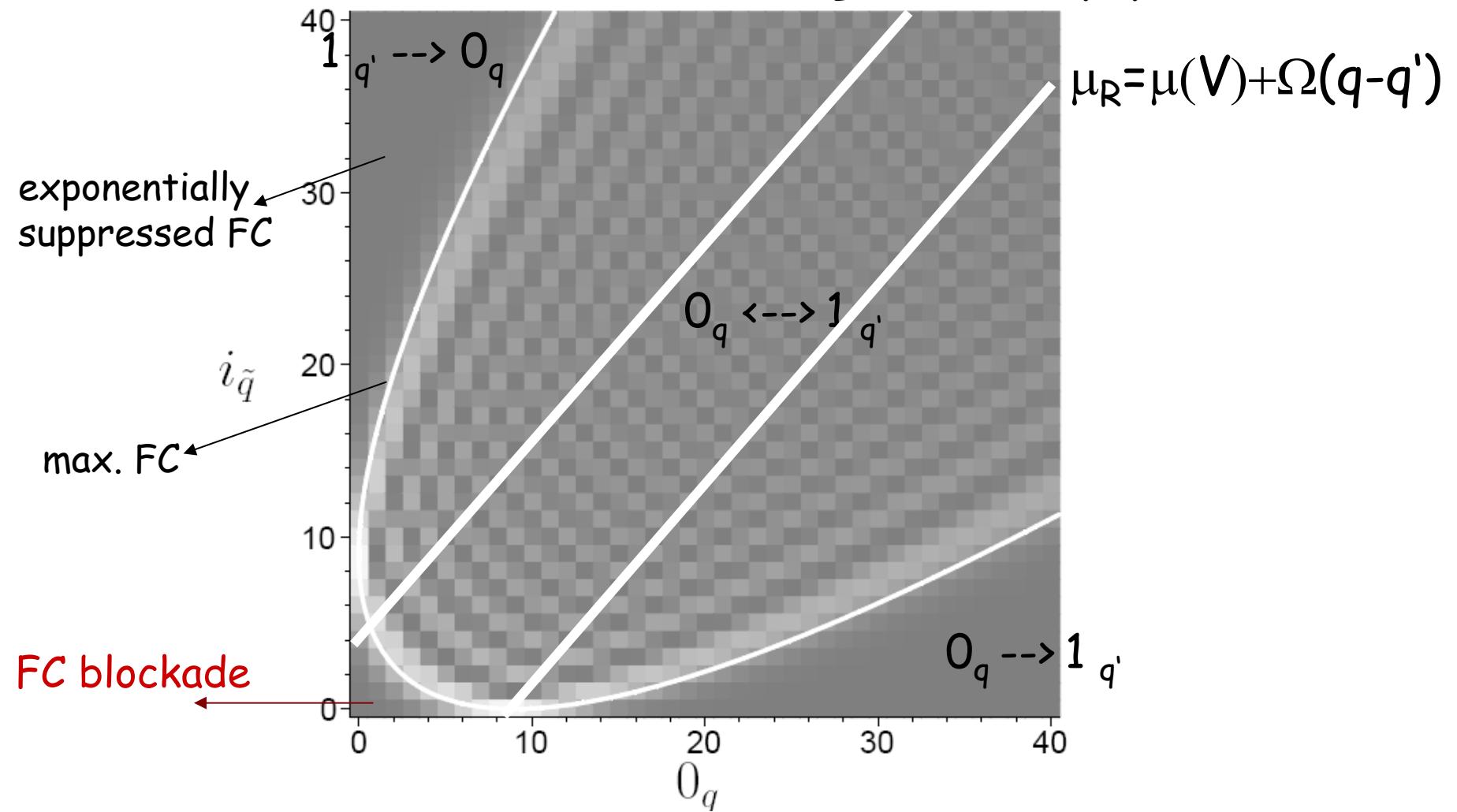
Weak charge-vibron coupling $g \ll 1$

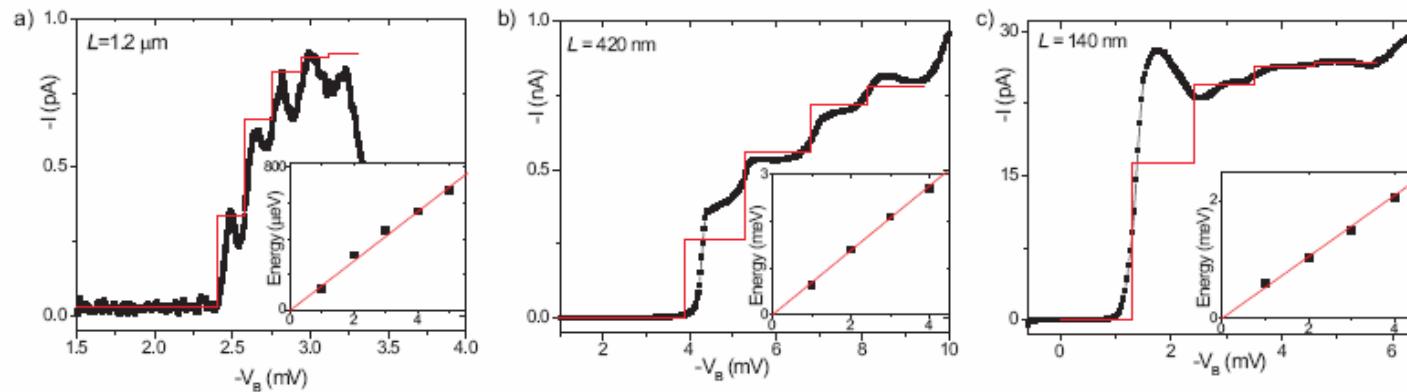


Strong charge-vibron coupling $g \gg 1$

$$\mu_L = \mu(V) + \Omega(q - q')$$

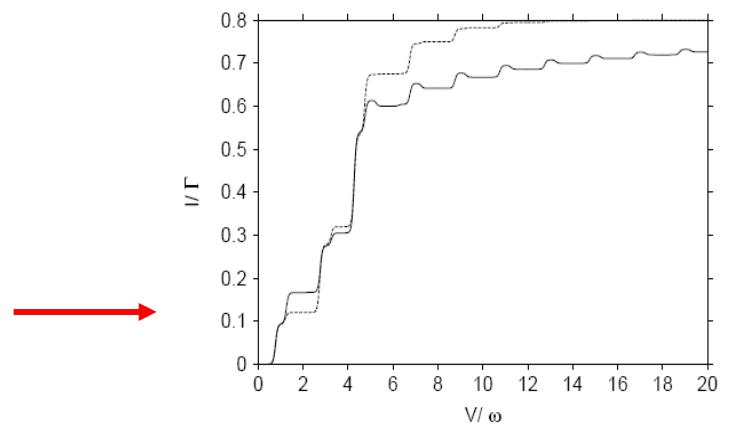
$$\mu_R = \mu(V) + \Omega(q - q')$$





Not clear: NDC alternating with PDC
Asymmetric vibron-charge coupling?
Nowack/Wegewijs, cond-mat/0506552

Further, fit is not perfect->
nonequilibrium phonons?
asymmetric coupling to leads?



Why electron-vibron coupling so large, $g \sim 1$?

In graphite g is small but
inhomogeneous electron charge distribution ?

Interaction with polarization charge:

$$W = - \int dx dx' \rho(x) K(x - x') \frac{\partial P}{\partial x'}. \quad P(x) \approx e \rho_0 z(x) \quad z(x) = A_n \sin(\pi n x / L)$$

→ Force $F = \frac{e \rho_0 \pi n}{L} \int_0^L dx \rho(x) \cos \frac{\pi n x}{L}. \quad \rightarrow \rho(x) = \text{const} \rightarrow F = 0 !$

Assume now $\rho(x) = e \delta(x - L/2)$ → $F_l = (-1)^l e^2 \rho_0 2\pi l / L$
only even harmonics $n=2l$

This is an extreme case, in general it is enough if $\rho(x)$ is not constant

Thanks to Andrea Donarini

The End