Charge Migration due to Electron Correlation in Molecules: Development of a Non-Dyson-method

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Overview

- Experimental Motivation What is Charge Migration?
- Dyson/CTA Method: Theory and Mechanisms of Charge Migration
- Development of the non-Dyson/CTA Method
- Illustrative Tests and Results of the non-Dyson/CTA

Experimental Motivation



Ionization by UV laser pulse ($\approx 14 \text{ eV}$). Initial hole localized at the benzene ring of amino acid tyrosin.

Charge Migration due to Electron Correlation in Molecules: Development of a Non-Dyson-method - p

Experimental Motivation



Bond near the N-terminal of the protein "on the left" breaks.

What is Charge Migration?

Charge migration =

Charge transfer solely driven by electron correlation

- Geometry of nuclei fixed (Born-Oppenheimer approximation)
- relevant timescale: few fs

Charge Transfer Analysis (CTA) developed in this group showed that efficient and ultrafast charge migration is, in fact, possible!



The Dyson/CTA method consists of

1. Dyson-ADC(3) scheme
 in order to gain
 the ionization spectrum
 3
 4
 4

ΔZ

Ú.

10

2. Subsequent Charge Transfer Analysis (CTA)

æ

30

ionization energy

30

- 1. What is Algebraic Diagrammatic Construction (ADC(3))?
 - Approximation of the many-particle Green's function via diagrammatic perturbation theory in 3rd order (Feynman graphs)
 - Terms of an algebraic expansion of a series are identified with terms resulting from Feynman graphs
- Dyson-ADC: Green's function determined via the Dyson equation Why do we use ADC?
 - Size consistency
 - Compactness
 - Experience with ADC, existence of programs

- 2. Charge transfer Analysis (CTA)
 - Ionization of neutral molecule in "exact" ground state $|\Psi_0\rangle$ (sudden approximation).
- => hole density

$$Q(\vec{r},t) := \underbrace{\langle \Psi_0 | \rho(\vec{r},t) | \Psi_0 \rangle}_{\rho_0(\vec{r})} - \underbrace{\langle \Phi_i | \rho(\vec{r},t) | \Phi_i \rangle}_{\rho_i(\vec{r},t)},$$

where

- ρ : charge density-operator
- $|\Phi_i\rangle$: cationic initial state

Charge density of cationic state:

$$\rho_{i}(\vec{r},t) = \sum_{I,J} \langle \Phi_{i} | I \rangle \langle I | e^{iHt} \rho(\vec{r},0) e^{-iHt} | J \rangle \langle J | \Phi_{i} \rangle$$

$$= \sum_{pq} \varphi_{p}^{*}(\vec{r}) \varphi_{q}(\vec{r}) \sum_{I,J} x_{I} \underbrace{\langle I | a_{p}^{\dagger} a_{q} | J \rangle}_{\text{perturb. exp.}} x_{J} \cos((E_{I} - E_{J})t)$$

$$\underbrace{N_{pq}(t)}$$

where

- $\{|I\rangle\}$ basis of eigenstates of the ionized molecule
- $x_I = \langle \Phi_i | I \rangle$ transition amplitudes

• $\{\varphi_p\}$ one-particle HF-orbital basis, $\rho(\vec{r},0) = \sum_{pq} \varphi_p^*(\vec{r}) \varphi_q(\vec{r}) a_p^{\dagger} a_q$

Calculate $\langle I | a_p^{\dagger} a_q | J \rangle$ via perturbation theory. Ansatz:

$$|I\rangle = \sum_{j} c_{j}^{(I)} a_{j} |\Phi_{0}\rangle + \sum_{a,k < l} c_{akl}^{(I)} a_{a}^{\dagger} a_{k} a_{l} |\Phi_{0}\rangle + \dots$$

 \Rightarrow Hole density

$$Q(\vec{r},t) = \sum_{p,q} \varphi_p^*(\vec{r}) \varphi_q(\vec{r}) N_{pq}(t) = \sum_p |\tilde{\varphi}_p(\vec{r},t)|^2 \tilde{n}_p(t)$$

Diagonalization

Eigenfunctions $\{\tilde{\varphi}_p\}$ called *natural charge orbitals* Eigenvalues $\{\tilde{n}_p(t)\}$ called *hole occupation numbers*

Brief summary:

A numerical Dyson/CTA-calculation consists of

- 1. Hartree-Fock calculation
- 2. Dyson-ADC(3)
- 3. Diagonalization of the ADC-matrix
- 4. Charge Transfer Analysis (CTA)

Mechanisms of Charge Migration



Overview

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Developing a non-Dyson/CTA Method

Why developing a new method?

- Biophysical relevant systems are too large for Dyson/CTA calculations
- Bottleneck is diagonalization of ADC-matrix

Idea: Circumvent diagonalization of ADC-matrix by propagating in the basis, in which the ADC-matrix is represented.

What is non-Dyson-ADC?

- Perturbation expansion directly of Green's function, not using Dyson equation \Rightarrow "non-Dyson"-ADC
- Nondiagonal representation of Green's function basis is basis of "Intermediate States" (Intermediate state representation, ISR)

Developing a non-Dyson/CTA Method

Hole density

$$Q(\vec{r},t) := \underbrace{\langle \Psi_0 | \rho(\vec{r},t) | \Psi_0 \rangle}_{\rho_0(\vec{r})} - \underbrace{\langle \Phi_i | \rho(\vec{r},t) | \Phi_i \rangle}_{\rho_i(\vec{r},t)},$$

 ρ : charge density operator, $|\Phi_i\rangle$: cationic initial state

• Insert $1 = \sum |\tilde{\Psi}_{\mathcal{M}}\rangle \langle \tilde{\Psi}_{\mathcal{M}}|$, where $\{|\tilde{\Psi}_{\mathcal{M}}\rangle\}$: Basis of Intermediate States

$$\rho_{i}(\vec{r},t) = \sum_{\mathcal{M},\mathcal{N}} \langle \Phi_{i} | e^{iHt} \tilde{\Psi}_{\mathcal{M}} \rangle \langle \tilde{\Psi}_{\mathcal{M}} | \rho(\vec{r},0) | \tilde{\Psi}_{\mathcal{N}} \rangle \langle \tilde{\Psi}_{\mathcal{N}} e^{-iHt} | \Phi_{i} \rangle$$

$$= \sum_{pq} \varphi_{p}^{*}(\vec{r}) \varphi_{q}(\vec{r}) \underbrace{\sum_{\mathcal{M},\mathcal{N}} \langle \Phi_{i}(t) | \tilde{\Psi}_{\mathcal{M}} \rangle \langle \tilde{\Psi}_{\mathcal{M}} | a_{p}^{\dagger} a_{q} | \tilde{\Psi}_{\mathcal{N}} \rangle \langle \tilde{\Psi}_{\mathcal{N}} | \Phi_{i}(t) \rangle}_{\tilde{N}_{pq}(t)}$$

Developing a non-Dyson/CTA Method

 \Rightarrow Hole density

$$Q(\vec{r},t) = \sum_{pq} \varphi_p^* \varphi_q \tilde{N}_{pq}(t),$$

$$\tilde{N}_{pq}(t) := \delta_{pq} n_p - \sum_{\mathcal{M}, \mathcal{N}} \langle \Phi_i(t) | \tilde{\Psi}_{\mathcal{M}} \rangle \underbrace{\langle \tilde{\Psi}_{\mathcal{M}} | a_p^{\dagger} a_q | \tilde{\Psi}_{\mathcal{N}} \rangle}_{\text{perturb. expansion}} \langle \tilde{\Psi}_{\mathcal{N}} | \Phi_i(t) \rangle$$

Main tasks

- Perturbation expansion of $\langle \tilde{\Psi}_{\mathcal{M}} | a_p^{\dagger} a_q | \tilde{\Psi}_{\mathcal{N}} \rangle$ in 2nd order ... and its "spinfree" formulation necessary for subsequent implementation
- Implementing a propagation $e^{-iHt} |\Phi_i\rangle = |\Phi_i(t)\rangle$

Analytical Comparison of the Methods

In which way do the two methods differ analytically?

The newly developed non-Dyson/CTA and the so far used Dyson/CTA

- ▲ build up on non-Dyson-ADC resp. Dyson-ADC ⇒ ionization spectra are slightly different
- are identical in terms of zeroth order

Do the numerical results of the two methods show conformity?

 \Rightarrow Propagate while using the "complete" eigenvalue spectrum (resulting from a full diagonalization of the non-Dyson-ADC matrix) and compare with corresponding Dyson/CTA results

Testing non-Dyson/CTA

Comparison of ionization spectra

- Molecule: FC₄NH₂
- Basis: DZP (double zeta basis with add. polarisation functions)
- Full diagonalization of both matrices
- Black: Dyson/CTA
 red: non-Dyson/CTA

 $F-C\equiv C-C\equiv C-NH_2$

Good agreement.

Testing non-Dyson/CTA

Comparison of time-dependent hole occupation numbers



Testing non-Dyson/CTA



Surprisingly good agreement!

Results of the non-Dyson/CTA



Consider ionization out of orbital 13a' (lines marked green)

Results of the non-Dyson/CTA



Results of the non-Dyson/CTA



Very efficient mechanism of charge migration over a peptid bond. Dinstinct ultrafast charge migration due to electron correlation.

Brief Summary

- An ab initio method called Dyson/CTA to investigate charge migration was presented. Problem: calculating larger systems
- A new non-Dyson/CTA was developed for a charge migration analysis towards investigating biophysical relevant systems
- Analytical comparison of the formulas of the so far used Dyson/CTA with the newly developed non-Dyson/CTA method
- Comparison of the numerical results of the new non-Dyson/CTA with Dyson/CTA and presentation of results of the non-Dyson/CTA
 using full diagonalization of non-Dyson-ADC matrix

Outlook

- Potential of non-Dyson/CTA method: Circumvent diagonalization of non-Dyson-ADC matrix (propagation routine implemented, tested with random matrices)
- Implementation of non-Dyson/CTA method represents a first step towards the analysis of charge migration in biologically relevant systems.

Further projects...

- Apply a complex absorbing potential (CAP) prepared
- Sudden approximation vs. shape of laser pulse

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