

Vibrational effects in the conductance through a molecular bridge

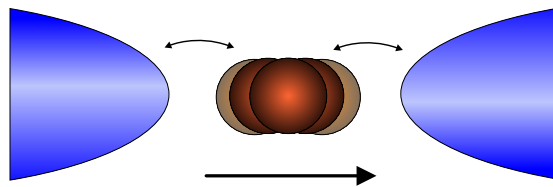
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(K. Richter, G. Cuniberti)

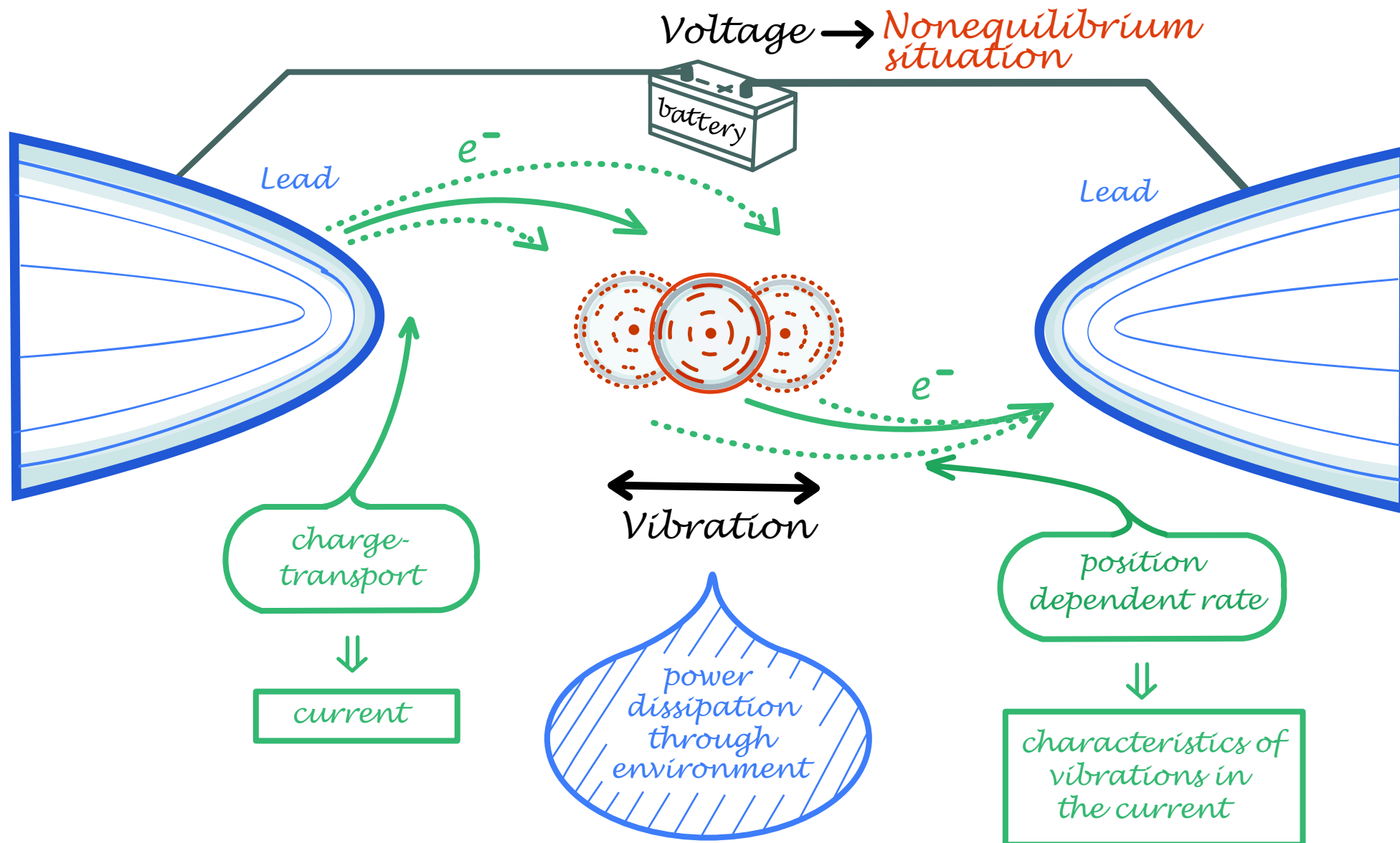
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Motivation

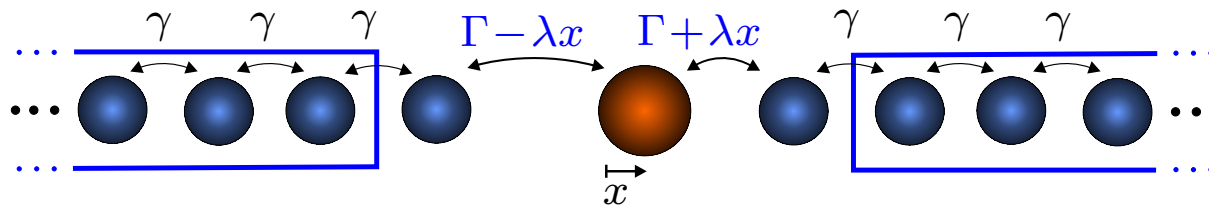
- standard calculations of quantum transport have reached a sophisticated level
- seldom the modelling of molecular vibrations are taken into account
- importance of vibrations demonstrated e.g. by J. van Ruitenbeek, Nature **419**, 906 (2002) ('Conductance through a single H_2 molecule')



Model



Model



$$H = H_{\text{MV}} + \hbar\omega_0(b^\dagger b + \frac{1}{2}) + H_{\text{leads}} + H_{\text{bath}}$$

central region coupled to a phonon mode:

$$H_{\text{MV}} = \begin{pmatrix} \epsilon_L & \Gamma & 0 \\ \Gamma & \epsilon_0 & \Gamma \\ 0 & \Gamma & \epsilon_R \end{pmatrix} + \begin{pmatrix} 0 & -\lambda x & 0 \\ -\lambda x & \lambda_0 x & \lambda x \\ 0 & \lambda x & 0 \end{pmatrix} \left. \begin{array}{l} \text{bond} \\ \text{onsite} \end{array} \right\} \text{coupling}$$

H_{leads} : semi-infinite linear chains

in a tight-binding description

and coupling to central region

H_{bath} : dissipative bath; set of harmonic oscillators

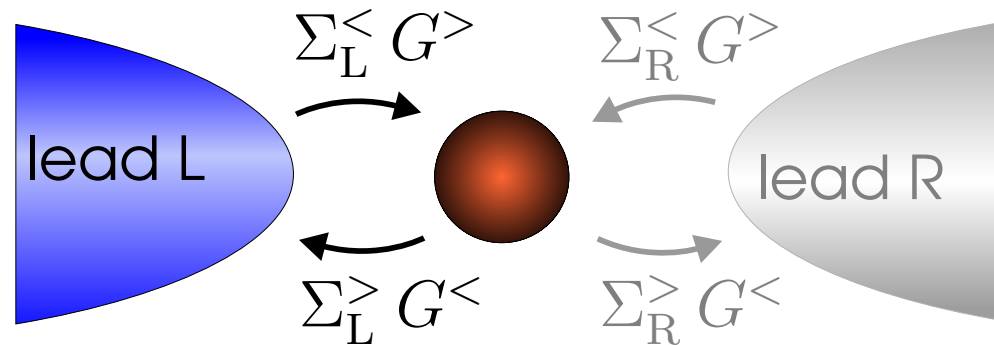
coupled to a phonon mode

Nonequilibrium current

total current: $I = \frac{1}{2}(I_L - I_R)$

$$I_\alpha = -2e \left\langle \frac{d}{dt} N_\alpha \right\rangle$$

$$= \frac{2e}{h} \int d\epsilon \left(\Sigma_\alpha^<(\epsilon) G^>(\epsilon) - \Sigma_\alpha^>(\epsilon) G^<(\epsilon) \right)$$



Meaning of greater and lesser Green functions:

$$G^<(\epsilon) = i f(\epsilon) A(\epsilon)$$

density of electrons

$$G^>(\epsilon) = -i(1 - f(\epsilon)) A(\epsilon)$$

density of holes

Dyson equation

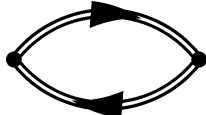
calculation of the electronic Green function:

$$\bar{G} = \bar{G}_0 + \bar{G}_0 (\bar{\Sigma}_L + \bar{\Sigma}_R + \bar{\Sigma}_{ph}) \bar{G}$$

with $\Sigma_{ph} =$  $+$ 

accounting for finite oscillation amplitudes:

$$\bar{D} = \bar{D}_0 + \bar{D}_0 (\bar{\Pi}_{el} + \bar{\Pi}_{bath}) \bar{D}$$

with $\Pi_{el} =$ 

$$\bar{G} = \begin{pmatrix} G^r & G^< + G^> \\ 0 & G^a \end{pmatrix}$$

Link to equilibrium theory

approximation for small voltages:

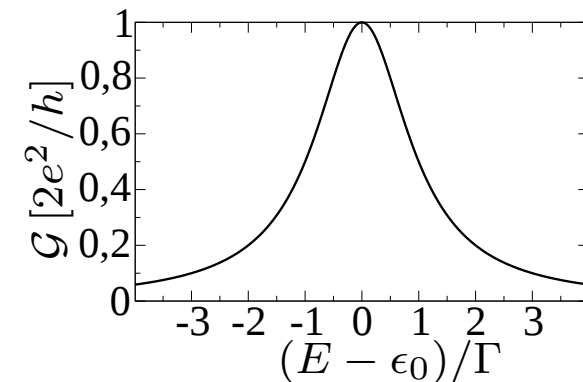
$$I = I_L - I_R = \frac{2e}{h} \int d\epsilon (f_L(\epsilon) - f_R(\epsilon)) \underbrace{\text{Tr}(G^a \Gamma_R G^r \Gamma_L)}_{\text{transmission}}$$

the zero-temperature conductance:

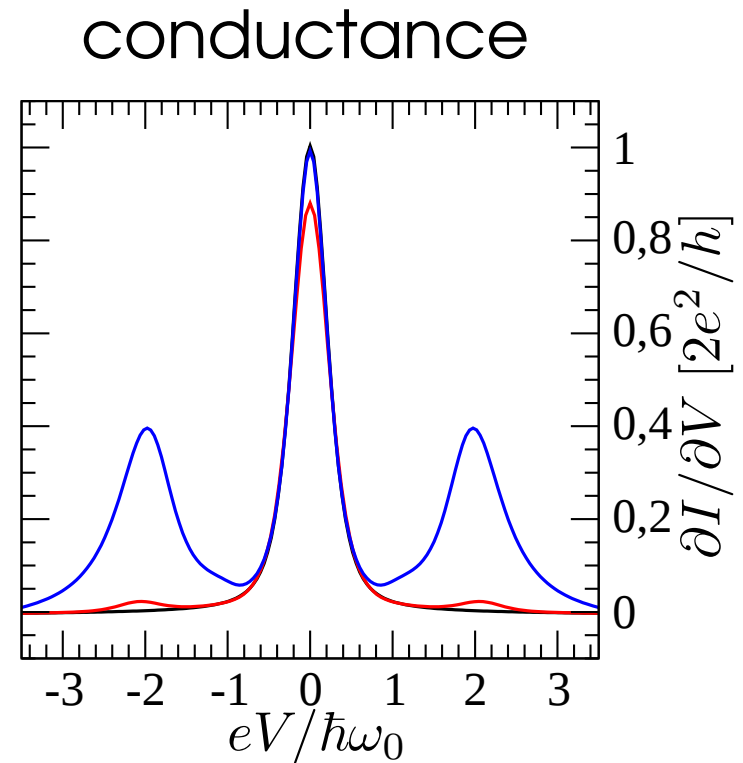
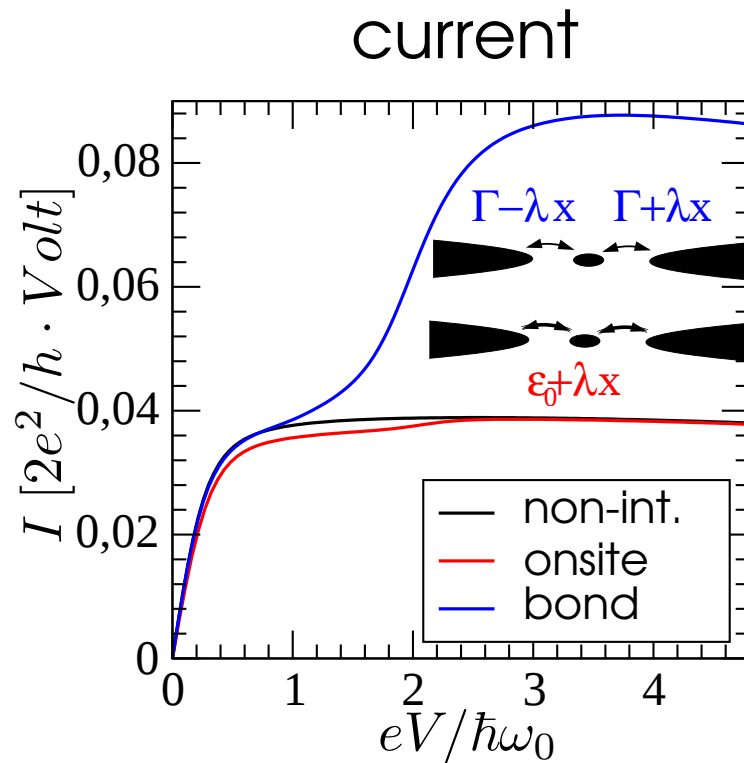
$$\mathcal{G} = \frac{\partial I}{\partial V} = \frac{2e^2}{h} T(\mu) = \frac{2e^2}{h} \text{Tr}(G^a \Gamma_R G^r \Gamma_L)$$

single energy level between two leads:

$$\mathcal{G} = \frac{2e^2}{h} \frac{\Gamma_L \Gamma_R}{(E - \epsilon_0)^2 + (\Gamma_L + \Gamma_R)^2/4}$$

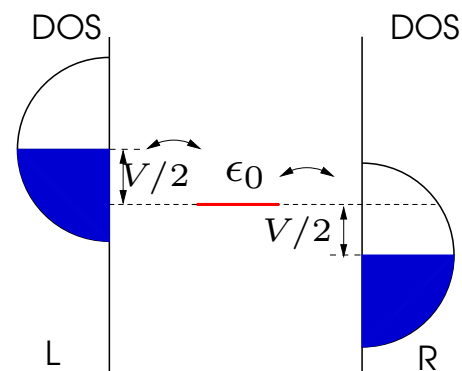


Results I: coupling mechanism

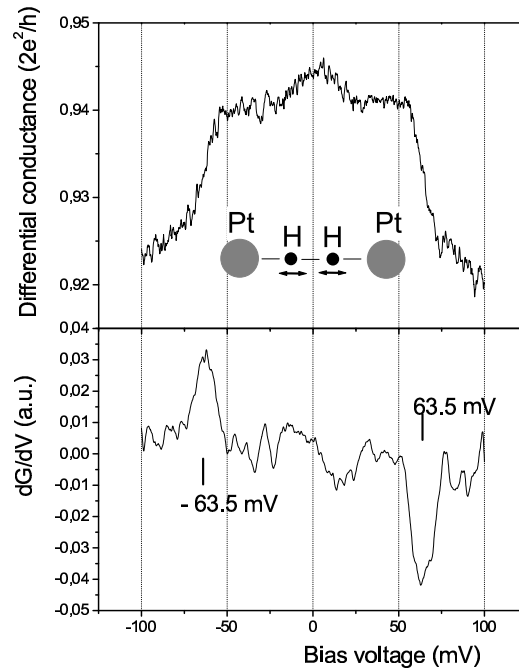


vibrational coupling:

$$\lambda = 0.24 \hbar\omega_0$$

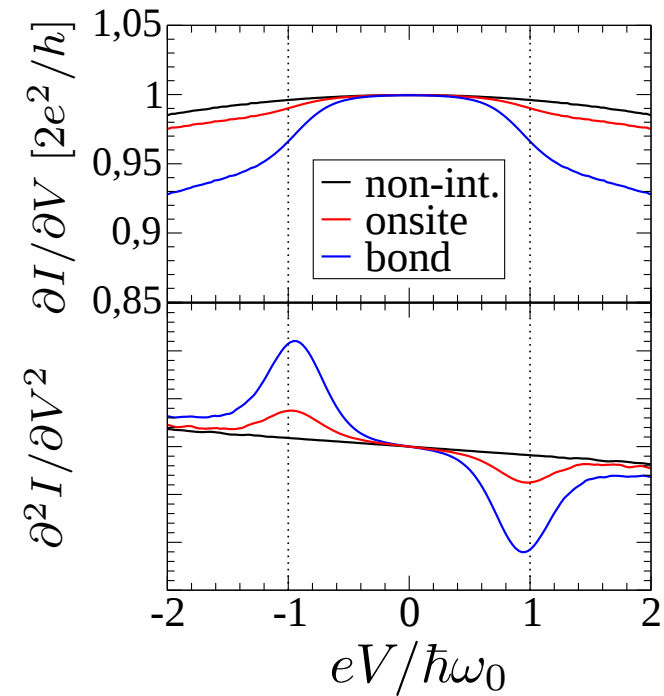


Results II: strong coupling regime



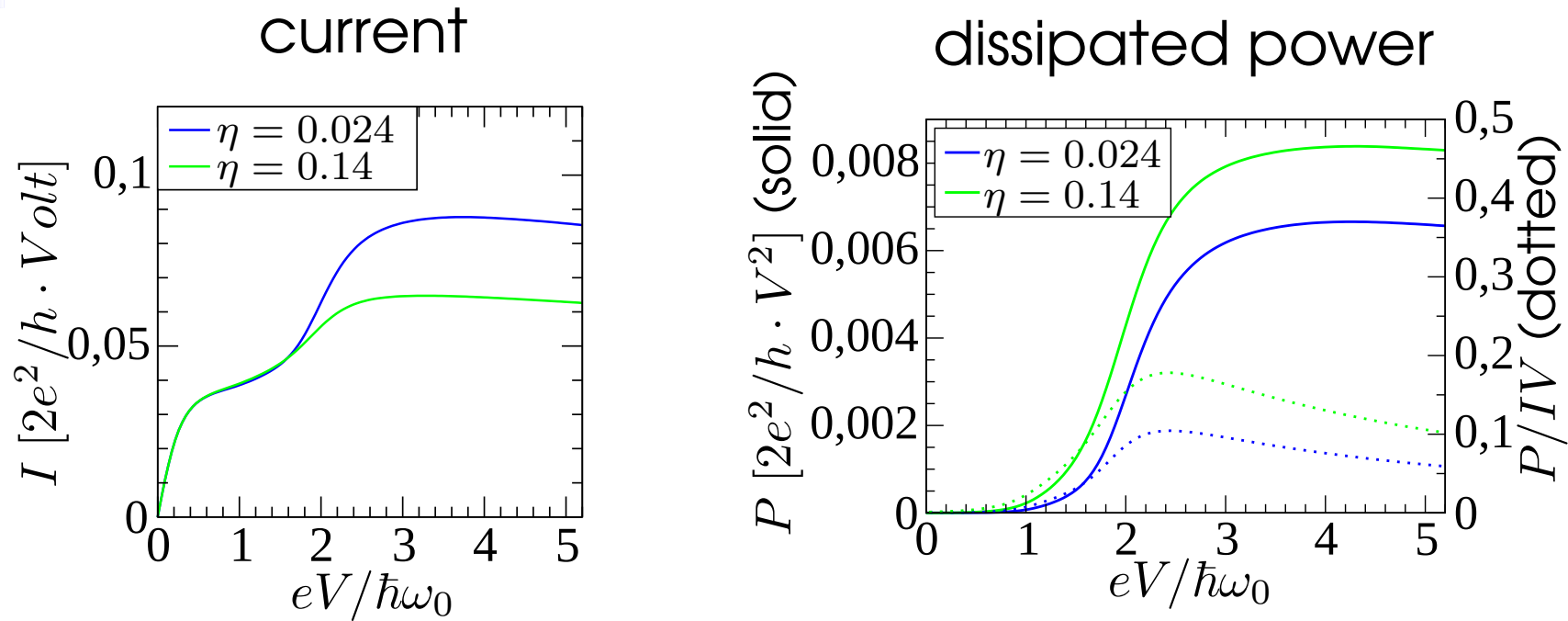
Conductance through a single H_2 -molecule.

Measurement by J. van Ruitenbeek, Nature **419**, 906 (2002)



Calculation of the center of mass vibration

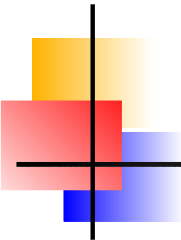
Results III: dissipated power



η : coupling strength to dissipative environment

- lead atoms
- other internal molecular vibrations
- surrounding

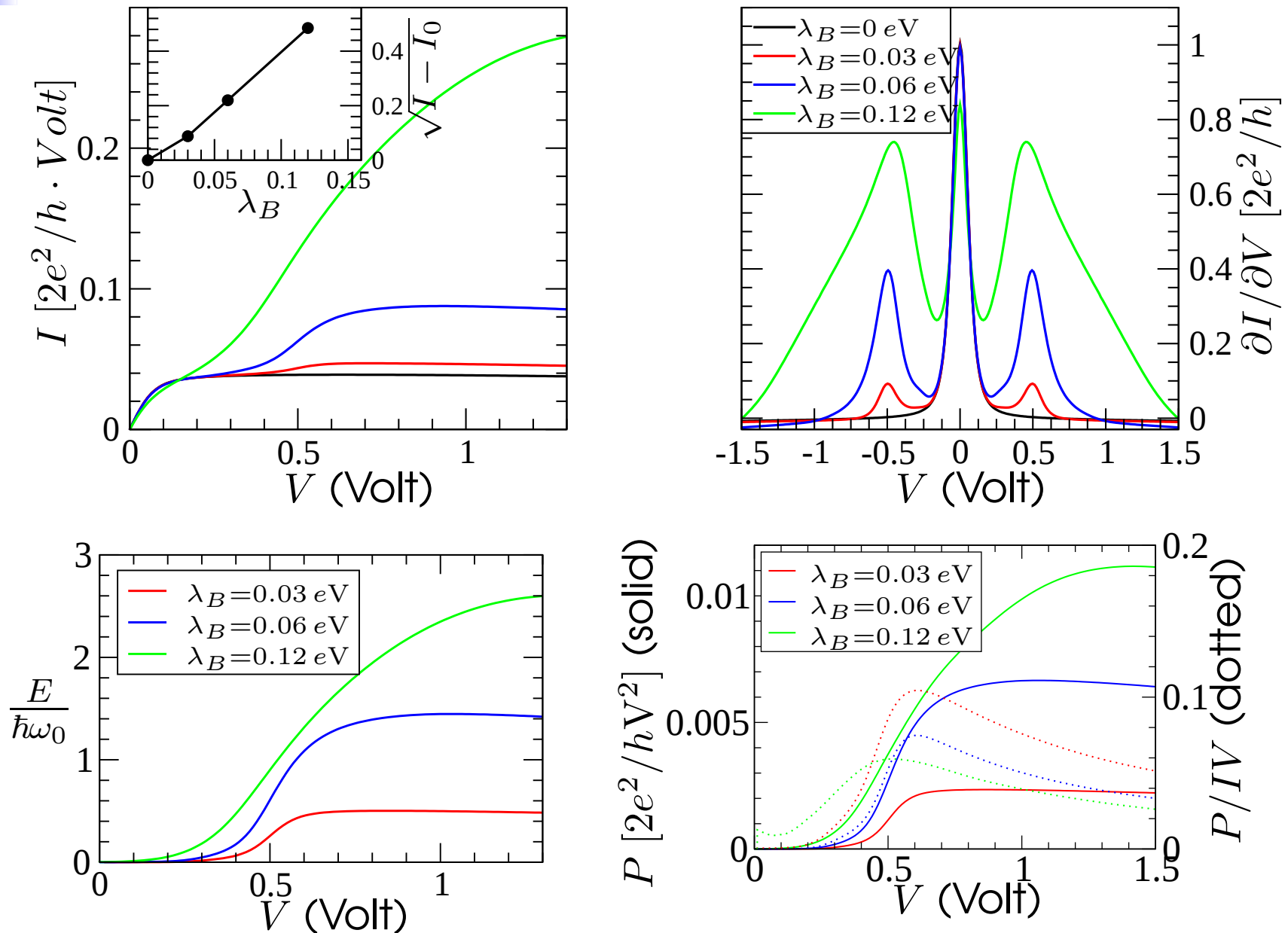
Modelling of bath: set of harmonic oscillators (ohmic)



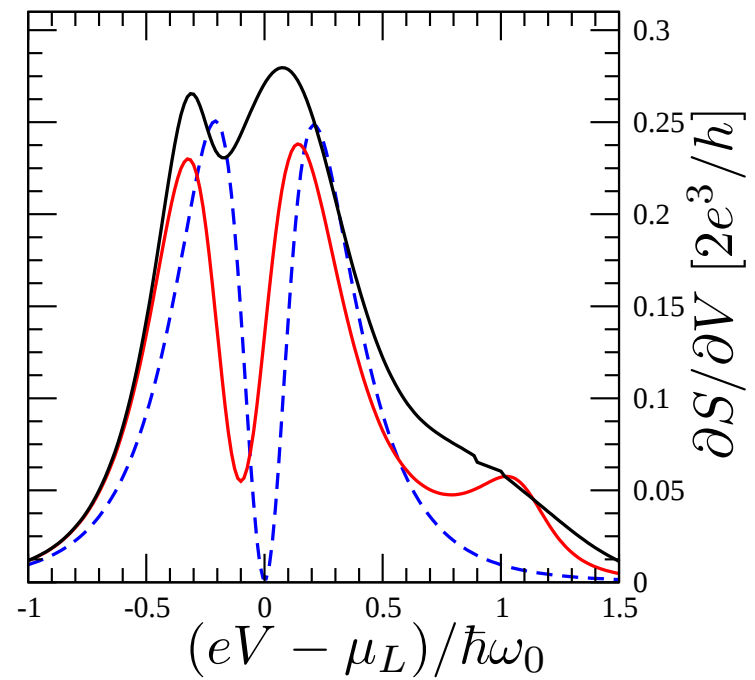
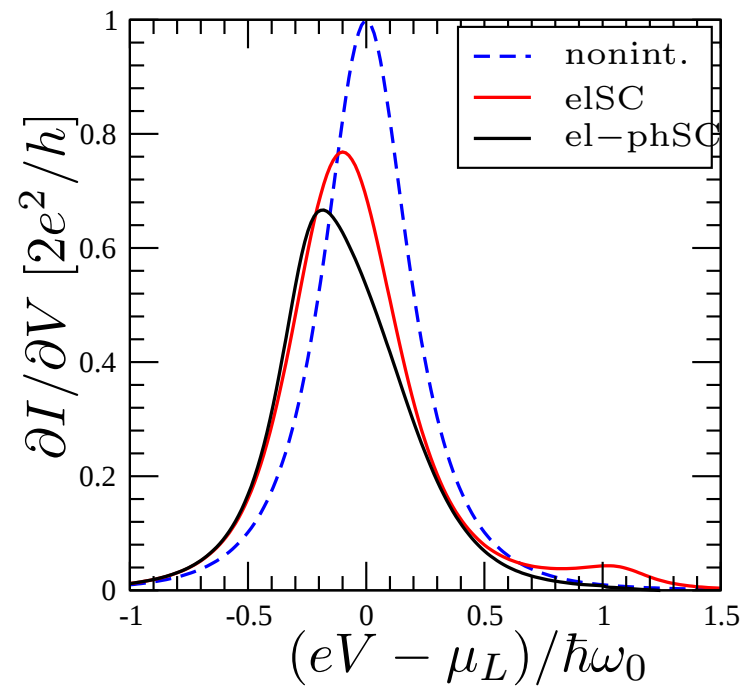
Conclusions and outlook

- ☀ relevance of bond stretching over onsite electron-phonon coupling
 - ☀ bond-stretching assisted current
 - ☀ quantum power $\neq$ classical power
-
- ? Transferability to DFT based calculations
 - ? Strong electron-phonon coupling
 - ? Coulomb blockade and Kondo regimes

Results IV: coupling strength



Results V: noise



Results VI: asymmetric

